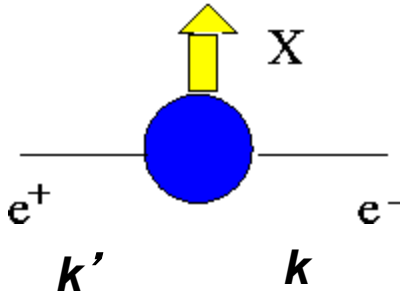


Riassunto della lezione precedente

- verifica di QPM in reazioni elettrodeboli, regole di somma :
 1. Gottfried sum rule \rightarrow $SU_f(3)$ rotta per i quark del mare
 2. Momentum sum rule \rightarrow gluoni portano metà del momento dell'adrone
- verifica fattorizzazione tra sez. d'urto elementare e densità partoniche
universalità di densità partoniche \rightarrow esplorare nuovi processi fondamentali
- Drell-Yan: cinematica, formule generali ; QPM picture \rightarrow test di N_c
- test sperimentali del QPM :
 - scaling di sezione d'urto
 - trend $\sigma(\pi^+A)/\sigma(\pi^-A)$ e $\sigma(\pi^-N)/\sigma(NN)$
 - distribuzione angolare della coppia leptonica $\sim(1+\cos^2\theta)$
 - \Leftrightarrow annichilazione elementare on-shell e collineare
- discrepanze :
 - K factor
 - distribuzione angolare $\neq (1+\cos^2\theta) \Rightarrow$ off-shell + radiation effects
 - meccanismo elementare diverso da QPM vicino a risonanze $\neq \gamma^*$

e^+e^- inclusivo (adronico)



$$q = k+k' \quad \text{time-like} \quad q^2 \equiv Q^2 = s \geq 0$$

$$d\sigma = \frac{1}{\mathcal{F}} |\mathcal{M}|^2 dR$$

$$\mathcal{F} = 4\sqrt{(k \cdot k')^2 - k^2 k'^2} \stackrel{\text{TRF}}{=} 2Q^2 \equiv 2s$$

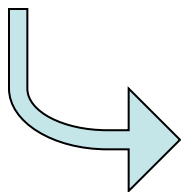
$$\mathcal{M} = \bar{v}(k') \gamma_\mu u(k) \frac{e^2}{Q^2} \langle P_X | J^\mu(0) | 0 \rangle \quad dR = (2\pi)^4 \delta(k + k' - P_X) \frac{d\mathbf{P}_X}{(2\pi)^3 2P_X^0}$$

$$|\mathcal{M}|^2 = \frac{e^4}{Q^4} L_{\mu\nu} H^{\mu\nu} \begin{cases} L_{\mu\nu} = 2(k_\mu k'_\nu + k'_\mu k_\nu - k \cdot k' g_{\mu\nu}) \\ H^{\mu\nu} = \sum_{S_X} \langle 0 | J^\mu | P_X \rangle \langle P_X | J^\nu | 0 \rangle \end{cases}$$

$$\sigma = \frac{1}{2} \frac{1}{2Q^2} \frac{e^4}{Q^4} L_{\mu\nu} \int \frac{d\mathbf{P}_X}{(2\pi)^3 2P_X^0} (2\pi)^4 \delta(q - P_X) H^{\mu\nu} = \frac{4\pi^2 \alpha^2}{Q^6} L_{\mu\nu} W^{\mu\nu}$$

media su polarizzazioni iniziali

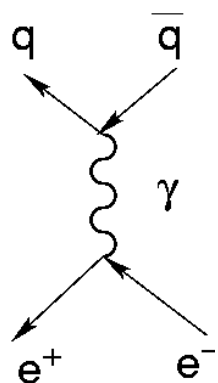
QPM picture



no adroni in stato iniziale e finale

σ in QPM $\equiv \sigma$ elementare $e^+e^- \rightarrow q\bar{q}$

solo N_c modi di creare la coppia conservando il colore nel vertice



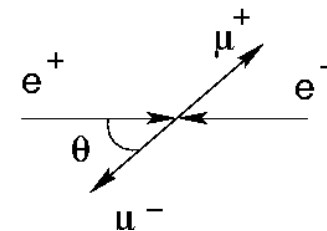
$Q^2 = s$ tale da avere solo produzione di γ

$$\sigma(e^+e^- \rightarrow q\bar{q}) \equiv \sigma(e^+e^- \rightarrow \mu^+\mu^-)$$

$$\sigma(e^+e^- \rightarrow X) = N_c \sum_f e_f^2 \sigma(e^+e^- \rightarrow q\bar{q})$$

$$= N_c \sum_f e_f^2 \int d\Omega \frac{d\sigma}{d\Omega}(e^+e^- \rightarrow \mu^+\mu^-)$$

$$= N_c \sum_f e_f^2 \int d\Omega \frac{\alpha^2}{4Q^2} (1 + \cos^2 \theta) = N_c \sum_f e_f^2 \frac{4\pi\alpha^2}{3Q^2}$$



Quindi

$$R = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = N_c \sum_f e_f^2$$

evidenza di N_c
 test di strutture
 $SU_c(3)$ e $SU_f(3)$

sotto soglia del c

$$R = 3 \left(\frac{4}{9} + \frac{2}{9} \right) = 2$$

vicino soglia

risonanze $J/\psi, \psi'$

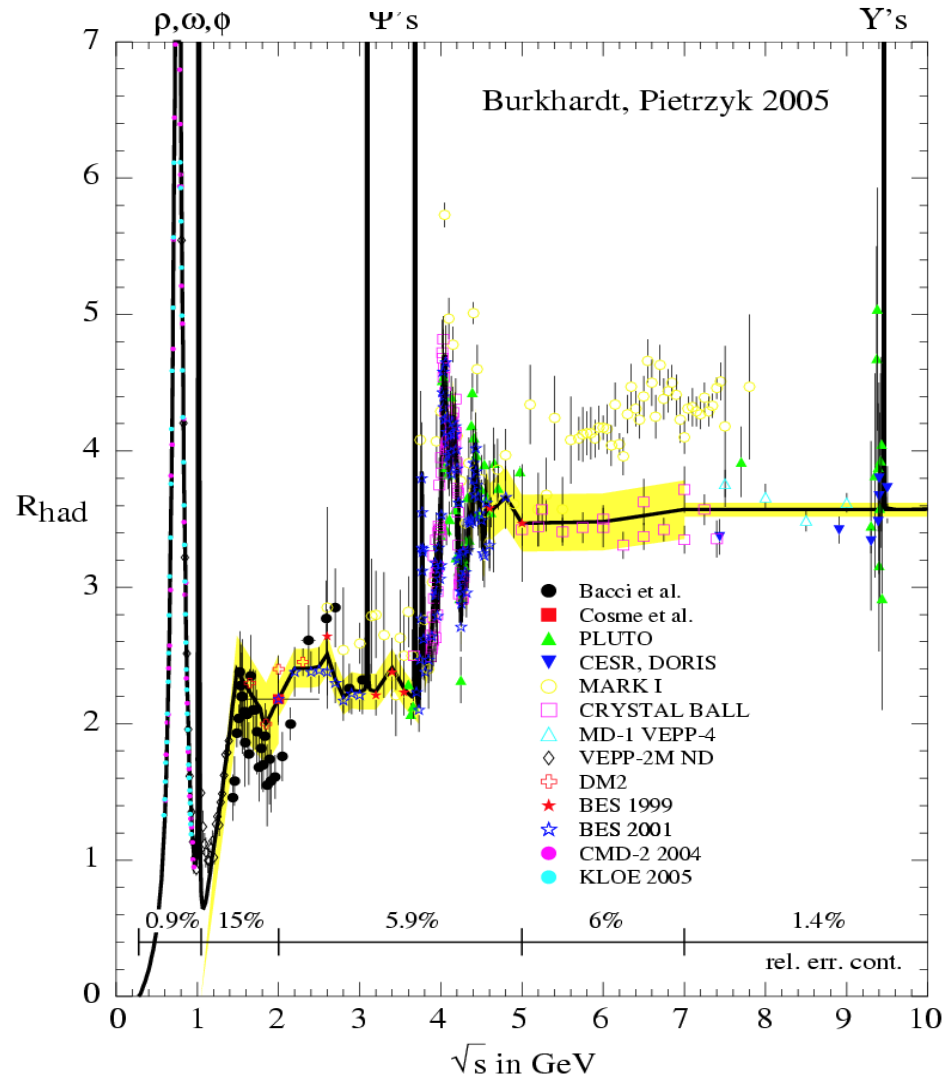
sopra soglia del c

$$R = 2 + 3 \frac{4}{9} = 3 + \frac{1}{3}$$

.....

oppure in

Wu, Phys.Rep. **C107** 59 (84)



e^+e^- semi-inclusivo adronico

crossing con DIS inclusivo

particelle ultrarelativistiche

$$k^\pm{}^\mu = (E, 0, 0, -/+ E)$$

$$q^\mu = (2E, \mathbf{0})$$

$$P_h^\mu = (E_h, E_h \sin \theta, 0, E_h \cos \theta)$$

invarianti

$$z = \frac{2P_h \cdot q}{q^2} \sim \frac{4EE_h}{4E^2} = \frac{E_h}{E}$$

$\sim "x_B^{-1}"$, misura elasticità

$$E_h \sim 0 \quad z \rightarrow 0$$

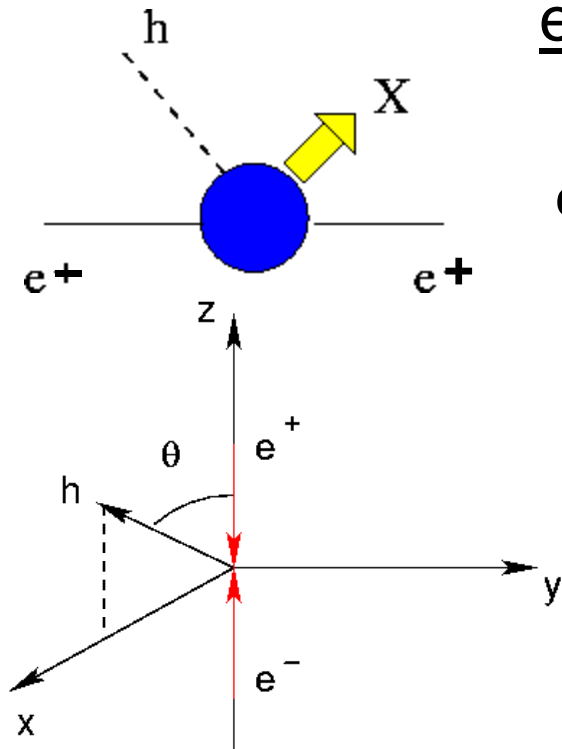
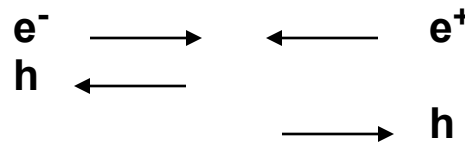
$$E_h \sim E \quad z \rightarrow 1$$

processo
elastico

$$y = \frac{P_h \cdot k'}{P_h \cdot q} \sim \frac{EE_h}{2EE_h} (1 + \cos \theta) = \frac{1}{2} (1 + \cos \theta) \sim \text{rapidity}$$

$$y = 0 \rightarrow \theta = \pi$$

$$y = 1 \rightarrow \theta = 0$$



e^+e^- semi-inclusivo adronico (continua)

$$d\sigma = \frac{1}{\mathcal{F}} |\mathcal{M}|^2 dR \quad \mathcal{F} = 4\sqrt{(k \cdot k')^2 - k^2 k'^2} \stackrel{\text{TRF}}{=} 2Q^2 \equiv 2s$$

$$dR = (2\pi)^4 \delta(q - P_X - P_h) \frac{d\mathbf{P}_X}{(2\pi)^3 2P_X^0} \frac{d\mathbf{P}_h}{(2\pi)^3 2E_h}$$

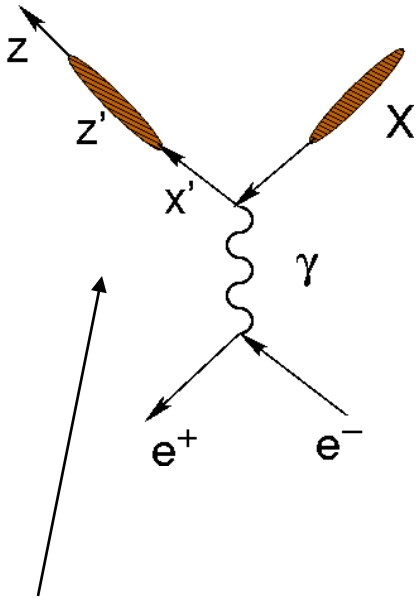
$$|\mathcal{M}|^2 = \frac{e^4}{Q^4} L_{\mu\nu} H_{1PI}^{\mu\nu}$$

$$\begin{aligned} 2E_h \frac{d\sigma}{d\mathbf{P}_h} &= \frac{1}{2} \frac{1}{2Q^2} \frac{1}{(2\pi)^2} \frac{e^4}{Q^4} L_{\mu\nu} \frac{1}{2\pi} \int \frac{d\mathbf{P}_X}{(2\pi)^3 2P_X^0} (2\pi)^4 \delta(q - P_X - P_h) H_{1PI}^{\mu\nu} \\ &= \frac{\alpha^2}{Q^6} L_{\mu\nu} W_{1PI}^{\mu\nu} \end{aligned}$$

N.B. cross-check $\sigma = \int \frac{d\mathbf{P}_h}{2E_h} 2E_h \frac{d\sigma}{d\mathbf{P}_h} = \frac{4\pi^2 \alpha^2}{Q^6} L_{\mu\nu} W^{\mu\nu}$



QPM picture



regime DIS : $Q^2 \rightarrow \infty$ con $z = \frac{2P_h \cdot q}{Q^2}$ finito

$$\begin{aligned} \frac{d\sigma}{dydz} &= N_c \sum_{f=q,\bar{q}} \int_0^1 dx' dz' \frac{d\sigma^{el}}{dx'dy} (e^+ e^- \rightarrow q\bar{q}) \\ &\quad \times D_f(z') \delta(z'x' - z) \\ &= N_c \sum_{f=q,\bar{q}} e_f^2 \int_0^1 dx' dz' \frac{d\sigma^{el}}{dy} (e^+ e^- \rightarrow \mu^+ \mu^-) \\ &\quad \times \delta(1 - x') D_f(z') \delta(z'x' - z) \end{aligned}$$

elastico \nearrow



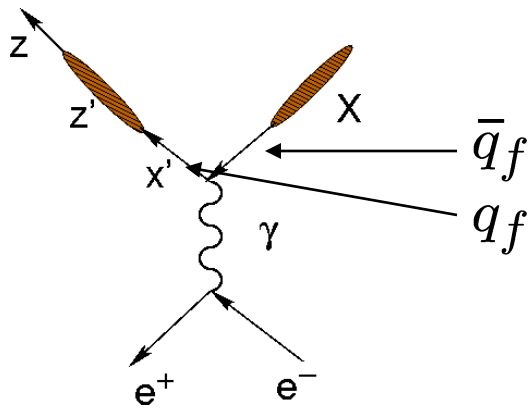
$$= N_c \frac{\pi\alpha^2}{Q^2} (1 + \cos^2 \theta) \sum_{f=q,\bar{q}} e_f^2 D_f(z)$$

probabilità di trovare adrone con frazione z' del momento x' del partone.

somma su tutte le combinazioni $x' z' \rightarrow$ adrone con frazione z dell' energia disponibile

commenti





conservazione del sapore nel vertice
 → stesso f per (anti)quark \Leftarrow Born approx.

correzioni ad ordini superiori da, e.g., g che decade in coppie → f' di antiquark $\neq f$ di quark

$$\frac{d\sigma}{dydz} = N_c \frac{\pi\alpha^2}{Q^2} (1 + \cos^2 \theta) \sum_{f=q,\bar{q}} e_f^2 D_f(z)$$

evidenza del colore

nuova incognita
 fattorizzazione
 → estrarre info da
 confronto con i dati

distribuzione angolare
 dell'adrone è data da
 una sez. d'urto di QED
 ($e^+e^- \rightarrow \mu^+\mu^-$) !

Fenomenologia e^+e^- semi-inclusivo

$$\sigma_{incl} = \sigma(e^+e^- \rightarrow X) = N_c \frac{4\pi\alpha^2}{3Q^2} \sum_f e_f^2$$

$$\sigma_{\mu\mu} = \int_0^1 dy \frac{d\sigma}{dy}(e^+e^- \rightarrow \mu^+\mu^-) = \frac{\pi\alpha^2}{Q^2} \int_0^1 dy (1 + \cos^2 \theta) = \frac{4\pi\alpha^2}{3Q^2}$$

$$\frac{d\sigma}{dz} = \int_0^1 dy \frac{d\sigma}{dzdy}(e^+e^- \rightarrow hX) = N_c \frac{4\pi\alpha^2}{3Q^2} \sum_{f=q,\bar{q}} e_f^2 D_f(z)$$

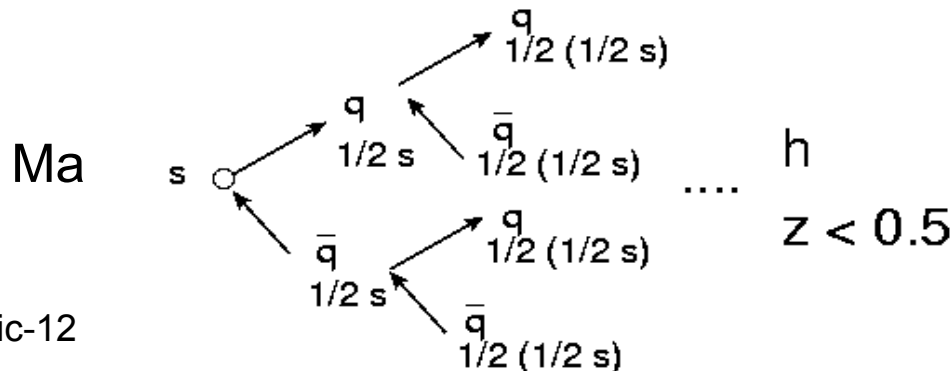


$$\frac{1}{\sigma_{\mu\mu}} \frac{d\sigma}{dz} = N_c \sum_f e_f^2 (D_f(z) + D_{\bar{f}}(z))$$

info su frammentazione

$$Q^2 \frac{d\sigma}{dz} = \frac{4\pi\alpha^2}{3} N_c \sum_f e_f^2 (D_f(z) + D_{\bar{f}}(z))$$

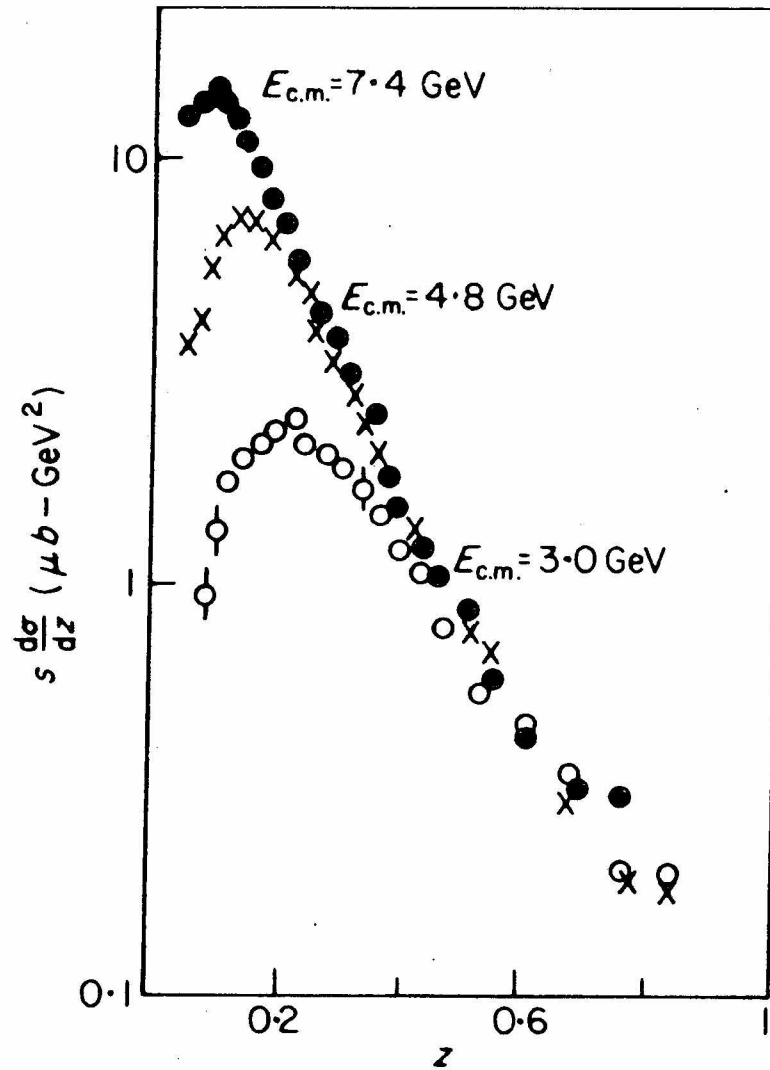
scaling in $z \forall Q^2 \equiv s$



violazione per $z \leq 0.5$



$$s \frac{d\sigma}{dz}(e^+e^- \rightarrow hX)$$



Schwitters *et al.*,
P.R.L. **35** 1320 (75)

FIG. 12.9. $s(d\sigma/dz)$ at 3, 4.8, and 7.4 GeV e^+e^- centre of mass energies.

e^+e^- semi-inclusivo \leftarrow fattorizzazione \rightarrow DIS ?

$$\frac{d\sigma}{dydz} = N_c \frac{\pi\alpha^2}{Q^2} (1 + \cos^2 \theta) \sum_f e_f^2 (D_f(z) + D_{\bar{f}}(z))$$

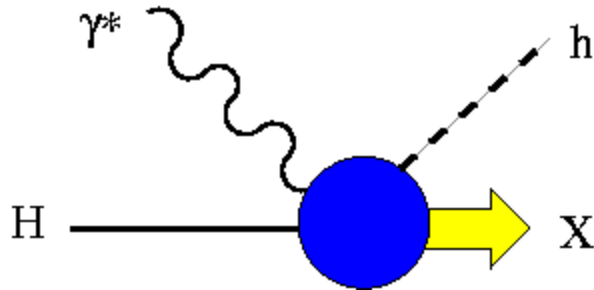


nuova incognita

fattorizzazione \rightarrow estrarre info da confronto con i
dati anche in altri processi

\Rightarrow DIS semi-inclusivo (SIDIS)

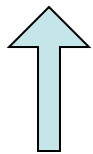
Semi-inclusive DIS (SIDIS)



stesse definizioni del caso inclusivo per cinematica e invarianti con in più

$$z_h = \frac{P \cdot P_h}{P \cdot q}$$

$z > 0$ P_h equiverso a $P \rightarrow h$ viene da frammentazione di partone
current fragmentation region



$z < 0$ P_h opposto a $P \rightarrow h$ viene da frammentazione del bersaglio H
target fragmentation region

N.B. in TRF $z \sim \frac{ME_h}{M\nu} = \frac{E_h}{\nu}$

SIDIS

$$\mathcal{F} = 4\sqrt{(P \cdot k)^2 - P^2 k^2} \stackrel{\text{TRF}}{=} 4ME \equiv 2s$$

$$d\sigma = \frac{1}{\mathcal{F}} |\mathcal{M}|^2 dR$$

$$dR = (2\pi)^4 \delta(k + P - k' - P_X - P_h) \frac{d\mathbf{P}_X}{(2\pi)^3 2P_X^0} \frac{d\mathbf{k}}{(2\pi)^3 2E'} \frac{d\mathbf{P}_h}{(2\pi)^3 2E_h}$$

$$|\mathcal{M}|^2 = \frac{e^4}{Q^4} L_{\mu\nu} H_{1PI}^{\mu\nu} \begin{cases} L_{\mu\nu} = 2(k_\mu k'_\nu + k'_\mu k_\nu - k \cdot k' g_{\mu\nu}) \\ H_{1PI}^{\mu\nu} = \frac{1}{2} \sum_{SS_h} \langle PS | J^\mu | P_X, P_h S_h \rangle \langle P_X, P_h S_h | J^\nu | PS \rangle \end{cases}$$

$$\begin{aligned} 2E_h \frac{d\sigma}{d\mathbf{P}_h dE' d\Omega} &= \frac{1}{4ME} \frac{E'}{16\pi^3} \frac{e^4}{Q^4} L_{\mu\nu} \frac{1}{(2\pi)^3} \int \frac{d\mathbf{P}_X}{(2\pi)^3 2P_X^0} (2\pi)^4 \delta(k + P - k' - P_X - P_h) H_{1PI}^{\mu\nu} \\ &= \frac{\alpha^2}{Q^4} \frac{E'}{E} L_{\mu\nu} \frac{1}{2M} 2M W_{1PI}^{\mu\nu} \end{aligned}$$



N.B. cross-check $\frac{d\sigma}{dE' d\Omega} = \int \frac{d\mathbf{P}_h}{2E_h} 2E_h \frac{d\sigma}{d\mathbf{P}_h dE' d\Omega} = \frac{\alpha^2}{Q^4} \frac{E'}{E} L_{\mu\nu} W^{\mu\nu}$

SIDIS (continua)

$$\frac{d\mathbf{k}'}{(2\pi)^3 2E'} = \frac{E' dE' d\Omega}{16\pi^3} = \frac{\pi y dx_B dy}{16\pi^3}$$

$$\frac{d\mathbf{P}_h}{2E_h} \sim d\mathbf{P}_{h\perp} \frac{dE_h}{2E_h} \sim d\mathbf{P}_{h\perp} \frac{dz}{2z}$$

$$\frac{d\sigma}{d\mathbf{P}_{h\perp} dx_B dy dz} = \frac{\pi\alpha^2}{Q^4} \frac{y}{2z} L_{\mu\nu} 2M W_{1PI}^{\mu\nu}$$



$$\frac{d\sigma}{dx_B dy dz} = \frac{\pi\alpha^2}{Q^4} \frac{y}{2z} L_{\mu\nu} 2M \int d\mathbf{P}_{h\perp} W_{1PI}^{\mu\nu}$$



cruciale per fattorizzazione tra
distribuzione e frammentazione

QPM picture

$$\begin{aligned} \frac{d\sigma}{dx_B dy dz} &= \sum_{f=q, \bar{q}} \int_0^1 dx' dz' \phi_f(x') \frac{d\sigma^{(e^- q \rightarrow e^- q)}}{dx' dy} \left(\frac{x'}{x_B} \right) D_f(z') \delta \left(z - z' \frac{x'}{x_B} \right) \\ &= \frac{4\pi\alpha^2 s}{Q^4} \left(\frac{y^2}{2} + 1 - y \right) x_B \sum_{f=q, \bar{q}} e_f^2 \phi_f(x_B) D_f(z) \end{aligned}$$



N.B. cross-check

$$\begin{aligned} \frac{d\sigma}{dx_B dy} &= \int dz \frac{d\sigma}{dx_B dy dz} \Big|_{D_f(z)=\delta(1-z)} \\ &= \frac{4\pi\alpha^2 s}{Q^4} \left(\frac{y^2}{2} + 1 - y \right) x_B \sum_{f=q, \bar{q}} e_f^2 \phi_f(x_B) = \frac{4\pi\alpha^2 s}{Q^4} \left(1 - y + \frac{y^2}{2} \right) F_2 \end{aligned}$$

Callan-Gross
 $F_2 = 2 x_B F_1$

DIS inclusivo

SIDIS

$$e^- p \rightarrow e'^- h^\pm X$$

$$\frac{1}{\sigma_{incl}} \frac{d\sigma}{dx_B dy dz} = \frac{\sum_f e_f^2 \phi_f(x_B) D_f(z)}{\sum_f e_f^2 \phi_f(x_B)} \sim D_u^h$$

e^+e^- semi-inclusivo

$$e^+ e^- \rightarrow h^\pm X$$

$$\frac{1}{\sigma_{incl}} \frac{d\sigma}{dz} = \frac{\sum_f e_f^2 D_f(z)}{\sum_f e_f^2} \sim D_u^h + D_{\bar{u}}^h$$

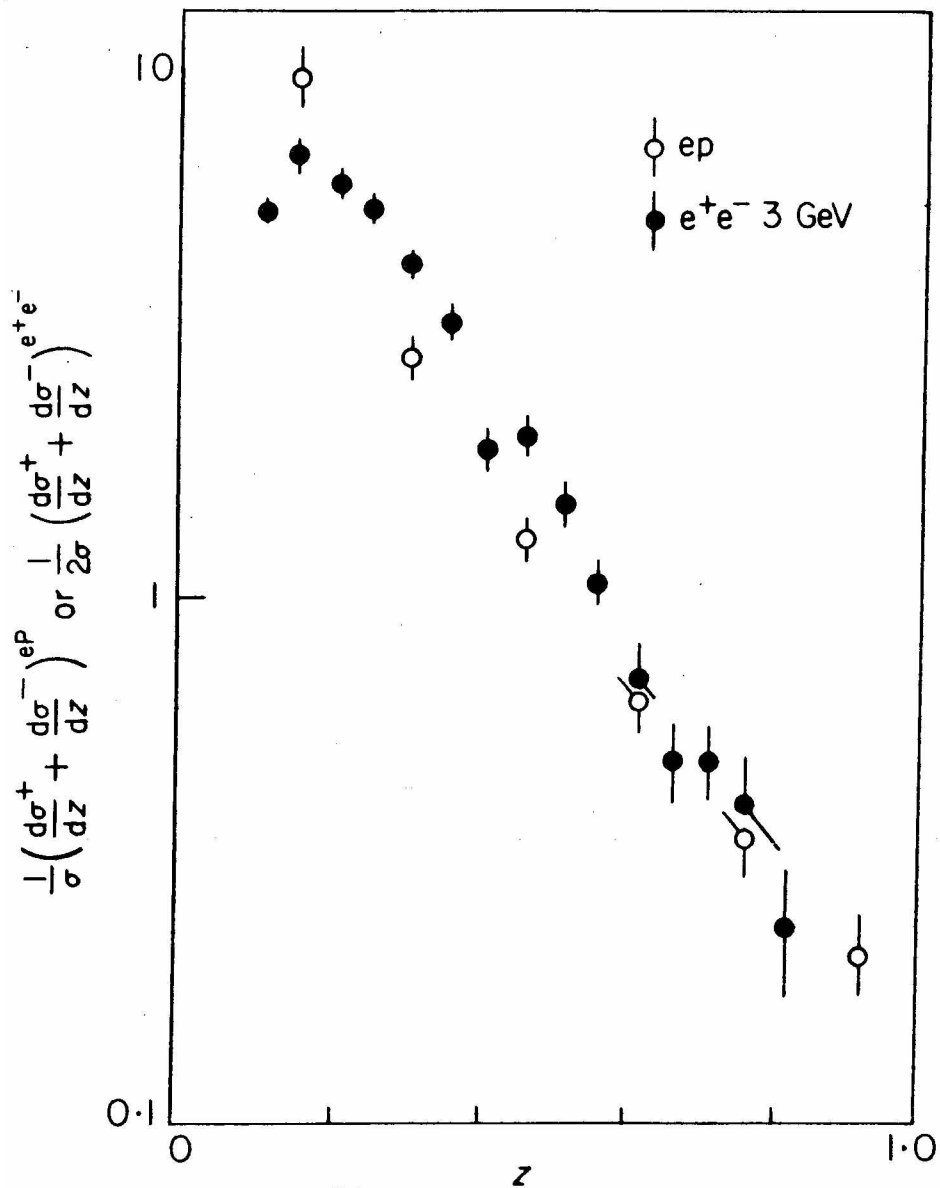
dominanza u quark

quindi

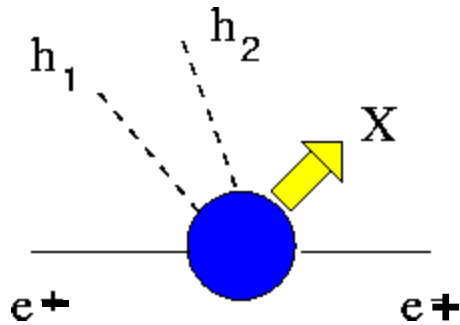


$$\frac{1}{\sigma_{incl}} \left(\frac{d\sigma^{h^+}}{dx_B dy dz} + \frac{d\sigma^{h^-}}{dx_B dy dz} \right) \Leftrightarrow \frac{1}{\sigma_{incl}} \frac{1}{2} \left(\frac{d\sigma^{h^+}}{dz} + \frac{d\sigma^{h^-}}{dz} \right) \sim D_u^{h^+} + D_{\bar{u}}^{h^-}$$

Gilman, Int. Symp. on lepton
and photon interactions
at high energies,
SLAC (75)



e^+e^- semi-inclusivo in due adroni



$$z_1 = \frac{2P_1 \cdot q}{q^2} \quad z_2 = \frac{2P_2 \cdot q}{q^2}$$

$$d\sigma = \frac{1}{\mathcal{F}} |\mathcal{M}|^2 dR \quad |\mathcal{M}|^2 = \frac{e^4}{Q^4} L_{\mu\nu} H_{2PI}^{\mu\nu}$$

$$\mathcal{F} \stackrel{\text{TRF}}{=} 2Q^2 \equiv 2s$$

$$dR = (2\pi)^4 \delta(q - P_X - P_1 - P_2) \frac{d\mathbf{P}_X}{(2\pi)^3 2P_X^0} \frac{d\mathbf{P}_1}{(2\pi)^3 2E_1} \frac{d\mathbf{P}_2}{(2\pi)^3 2E_2}$$

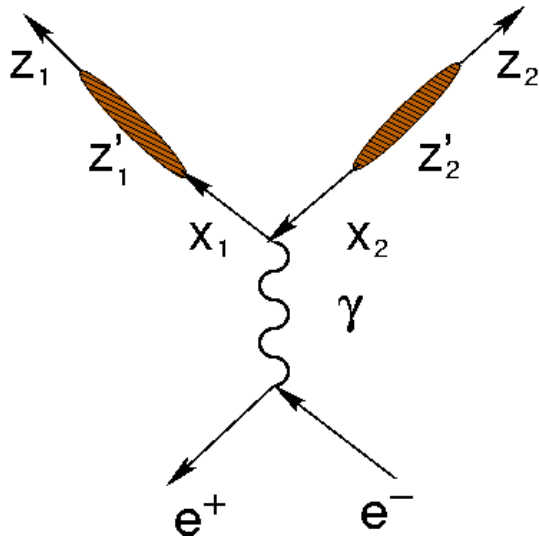
$$\begin{aligned} 2E_1 2E_2 \frac{d\sigma}{d\mathbf{P}_1 d\mathbf{P}_2} &= \frac{1}{2} \frac{1}{2Q^2} \frac{1}{(2\pi)^2} \frac{e^4}{Q^4} L_{\mu\nu} \\ &\quad \times \frac{1}{(2\pi)^4} \int \frac{d\mathbf{P}_X}{(2\pi)^3 2P_X^0} (2\pi)^4 \delta(q - P_X - P_1 - P_2) H_{2PI}^{\mu\nu} \\ &= \frac{\alpha^2}{Q^6} L_{\mu\nu} W_{2PI}^{\mu\nu} \end{aligned}$$



N.B. cross-check $\int \frac{d\mathbf{P}_2}{2E_2} 2E_1 2E_2 \frac{d\sigma}{d\mathbf{P}_1 d\mathbf{P}_2} = 2E_1 \frac{d\sigma}{d\mathbf{P}_1} = \dots = \frac{\alpha^2}{Q^6} L_{\mu\nu} W_{1PI}^{\mu\nu}$

QPM picture

$$\begin{aligned}
 \frac{d\sigma}{dydz_1dz_2} &= N_c \sum_{f=q,\bar{q}} \int_0^1 dx'_1 dx'_2 dz'_1 dz'_2 \frac{d\sigma^{el}}{dx'_1 dx'_2 dy} (e^+ e^- \rightarrow q\bar{q}) \\
 &\quad \times D_f(z'_1) \delta(z_1 - z'_1 x'_1) D_f(z'_2) \delta(z_2 - z'_2 x'_2) \\
 &= N_c \sum_f e_f^2 \int_0^1 dx'_1 dx'_2 dz'_1 dz'_2 \frac{d\sigma^{el}}{dy} (e^+ e^- \rightarrow \mu^+ \mu^-) \\
 &\quad \times \delta(1 - x'_1) D_f(z'_1) \delta(z_1 - z'_1 x'_1) \\
 &\quad \times \delta(1 - x'_2) D_f(z'_2) \delta(z_2 - z'_2 x'_2) \\
 &= N_c \frac{\pi\alpha^2}{Q^2} (1 + \cos^2 \theta) \sum_{f=q,\bar{q}} e_f^2 D_f(z_1) D_f(z_2)
 \end{aligned}$$



N.B. cross-check $D_f(z_2) = \delta(1 - z_2)$ secondo adrone \equiv jet adronico

$$\int dz_2 \frac{d\sigma}{dydz_1dz_2} \Big|_{D_f(z_2)=\delta(1-z_2)} = \frac{d\sigma}{dydz} = N_c \frac{\pi\alpha^2}{Q^2} (1 + \cos^2 \theta) \sum_f e_f^2 D_f(z)$$

e⁺e⁻ semi-inclusivo

Adesso $D_f(z_1) = \delta(1 - z_1)$ anche primo adrone \equiv jet adronico

$$\int dz \frac{d\sigma}{dydz} \Big|_{D_f(z)=\delta(1-z)} = \frac{d\sigma^{jet}}{dy} = N_c \frac{\pi\alpha^2}{Q^2} (1 + \cos^2 \theta) \sum_f e_f^2$$

distribuzione angolare di tutti gli adroni
nello stato finale
=
sezione d'urto di jet

$$\frac{d\sigma^{el}}{dy}(e^+e^- \rightarrow q\bar{q}) = \frac{\pi\alpha^2}{Q^2} \underline{(1 + \cos^2 \theta)} e_f^2$$

$$\frac{d\sigma}{dydz}(e^+e^- \rightarrow hX) = N_c \frac{\pi\alpha^2}{Q^2} \underline{(1 + \cos^2 \theta)} \sum_f e_f^2 D_f(z)$$

$$\frac{d\sigma}{dy}(e^+e^- \rightarrow \text{jets}) = N_c \frac{\pi\alpha^2}{Q^2} \underline{(1 + \cos^2 \theta)} \sum_f e_f^2$$

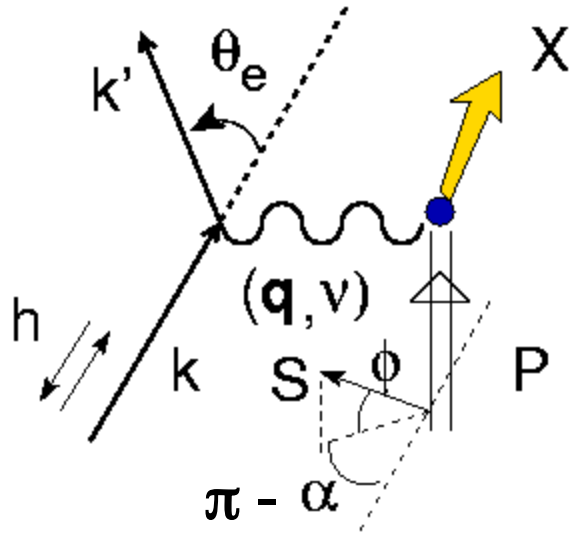
gli adroni sono “frammenti” dei partoni a spin 1/2 del processo elementare
 eventi a molti adroni = gruppi di adroni con p_T limitato rispetto ad un certo asse

dato asse θ , sfericità $S = \frac{3 \sum_i p_{Ti}^2}{2 \sum_i p_i^2} \begin{cases} S = 1 & \text{sfera} \\ S = 0 & \text{jet} \end{cases} \quad S \xrightarrow{s \rightarrow \infty} 0$

adrone in stato finale con $0 \leq z \leq 1$ si muove in un jet che rappresenta la direzione θ del quark di frammentazione rispetto all'asse z

la direzione del jet è data da processo elementare di QED !

DIS inclusivo polarizzato



se $S=0 \rightarrow$ violazione della parità
 processo debole \rightarrow corrente $V-A \rightarrow W^{\mu\nu}_A$

se $S \neq 0 \rightarrow$ 2 4-vettori P, q e
 1 4-pseudovettore S indipendenti
 struttura del tensore adronico più ricca

si sceglie S^μ tale che $S^2 = -1$ e $S \cdot P = 0$

$$S^\mu = \frac{S \cdot q}{P \cdot q} \left(P^\mu - \frac{M^2}{P \cdot q} q^\mu \right) + S^\mu_\perp = \frac{\lambda}{M} \left(P^\mu - \frac{M^2}{P \cdot q} q^\mu \right) + S^\mu_\perp$$

elicità

$$\lambda = M \frac{S \cdot q}{P \cdot q}$$

$$S \cdot P = 0 \rightarrow S_\perp \cdot P = 0$$

$$S^2 \sim -(\lambda^2 + S_\perp^2) = -1$$



Tensore adronico

$S = \frac{1}{2} \rightarrow W^{\mu\nu}$ è al più lineare in S , perchè è matrice 2×2 in spazio di spin

\Rightarrow espansione sulla base delle matrici σ di Dirac

$$W^{\mu\nu} = \sum_{\alpha\alpha'} W_{\alpha\alpha'}^{\mu\nu} \rho_{\alpha\alpha'} = \frac{1}{2} \sum_{\alpha\alpha'} W_{\alpha\alpha'}^{\mu\nu} (1 + \mathbf{P} \cdot \boldsymbol{\sigma})_{\alpha\alpha'}$$

matrice densità di spin del target

vettore di polarizzazione

$$P_i = \frac{N_+ - N_-}{N_+ + N_-} = \langle \sigma_i \rangle = \text{Tr}(\rho \sigma_i)$$

- S^u coplanar with scattering plane $\rightarrow \phi = 0$
- hermiticity del tensore
- invarianza per trasformazioni di parità
- invarianza per trasformazioni di time-reversal
- conservazione della corrente

