# Review of Semiconductor Drift Detectors

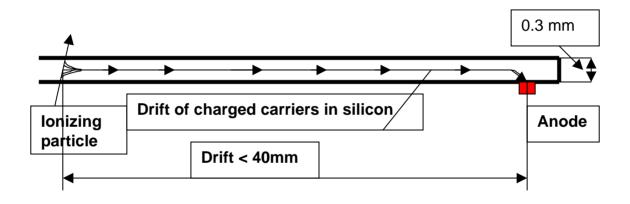
Talk given by Pavel Rehak following a presentation on 5<sup>th</sup> Hiroshima Symposium of Semiconductor Tracking Detectors

### Outline of the Review

- Principles of Semiconductor Drift Detectors
- Drift Detectors for Position and Charge sensing
- Lessons learned from utilization of Drift Detectors in Heavy Ion Experiments
- X-ray spectroscopy with Drift Detectors
- Controlled Drift Detectors
- Conclusions and predictions

#### Concept of Semiconductor Drift Detector

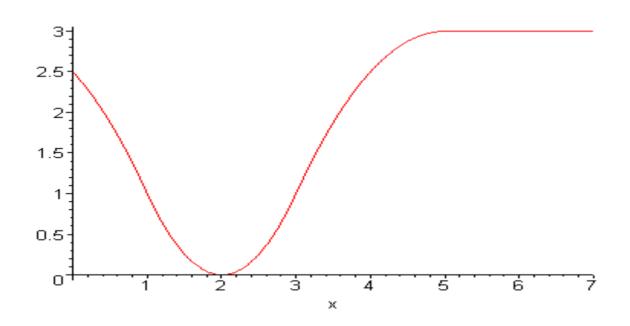
Transport of charged carriers in thin fully depleted semiconductor detectors in direction parallel to the large surface of the detector.



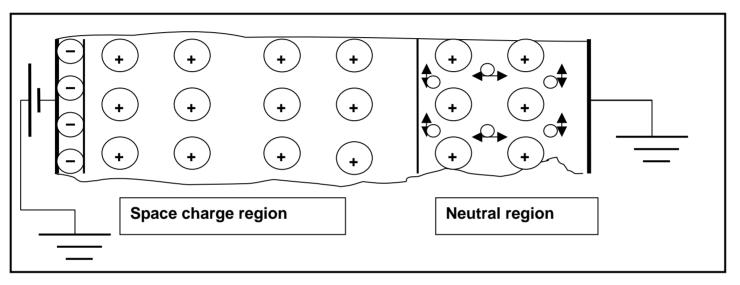
- 1.) Charge carriers induce a signal on the anode only on the arrival. The drift time measure the distance between the position of the ionizing particle and the anode.
- 2.) The capacitance of the anode is small -> low series noise of the read-out.

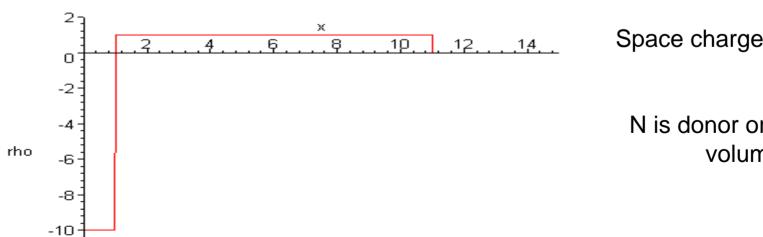
# Analyses of buried channel Charge Coupled Devices. (CCDs) in 1983

Negative potential in the region of the channel.



# Reversed biased one sided step p-n junction in 1dimension



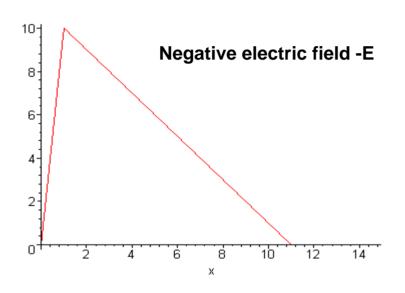


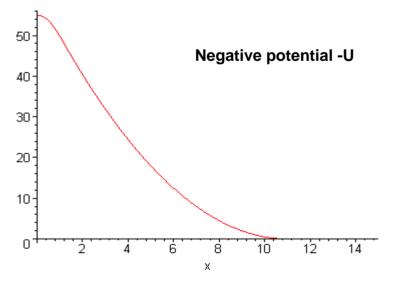
Space charge density ρ

 $\rho = qN$ 

N is donor or acceptor volume density

# Equations of electrostatics





div  $E=\rho/\varepsilon$ 

In one dimensional case the above reduces into:  $dE/dx = \rho/\epsilon$ 

Electric field is a piecewise linear.

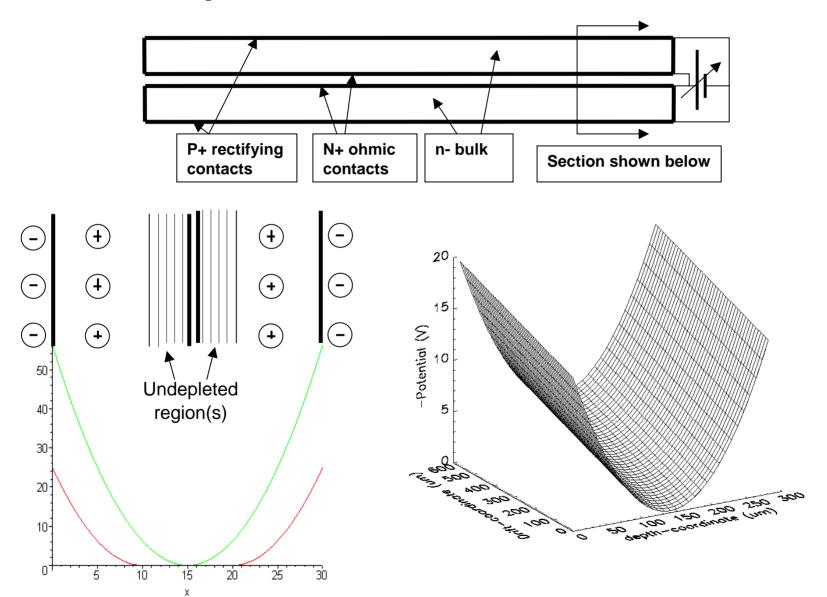
Electric potential *U* is defined by:

E=-grad U

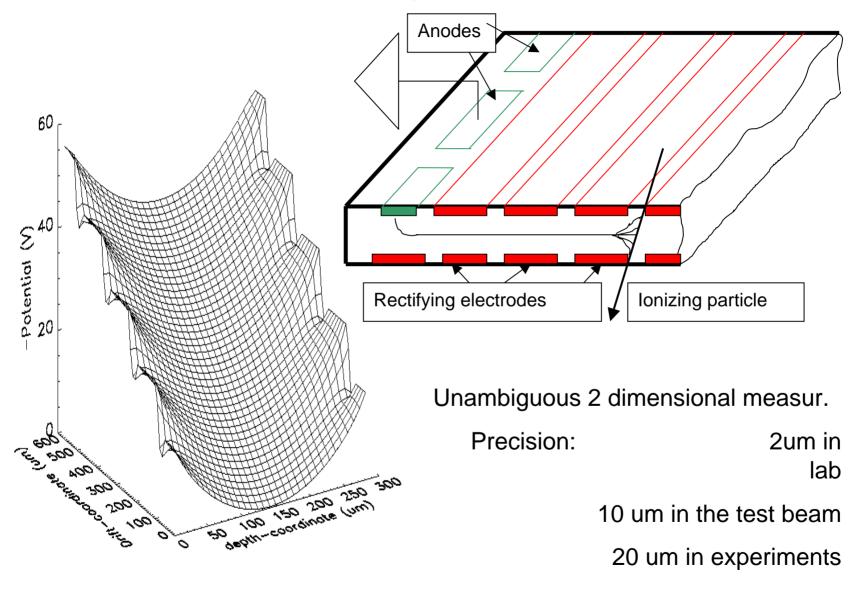
In one dimensional case the above reduces into: *E=-dU/dx* 

U is parabolic and depleted depth d  $d = \sqrt{2\varepsilon U/qN}$ 

## Depletion from two sides



# Summary of principles



# Optimal signal processing

$$\varepsilon^{2} = \frac{C^{2}e_{n}^{2}/2\int_{-\infty}^{+\infty}w^{2}(t)dt + q^{2}v\int_{-\infty}^{+\infty}w^{2}(t)dt + q^{2}N\int_{-\infty}^{+\infty}f(t)w^{2}(t)dt}{q^{2}N^{2}\left[\int_{-\infty}^{+\infty}f'(t)w(t)dt\right]^{2}}$$

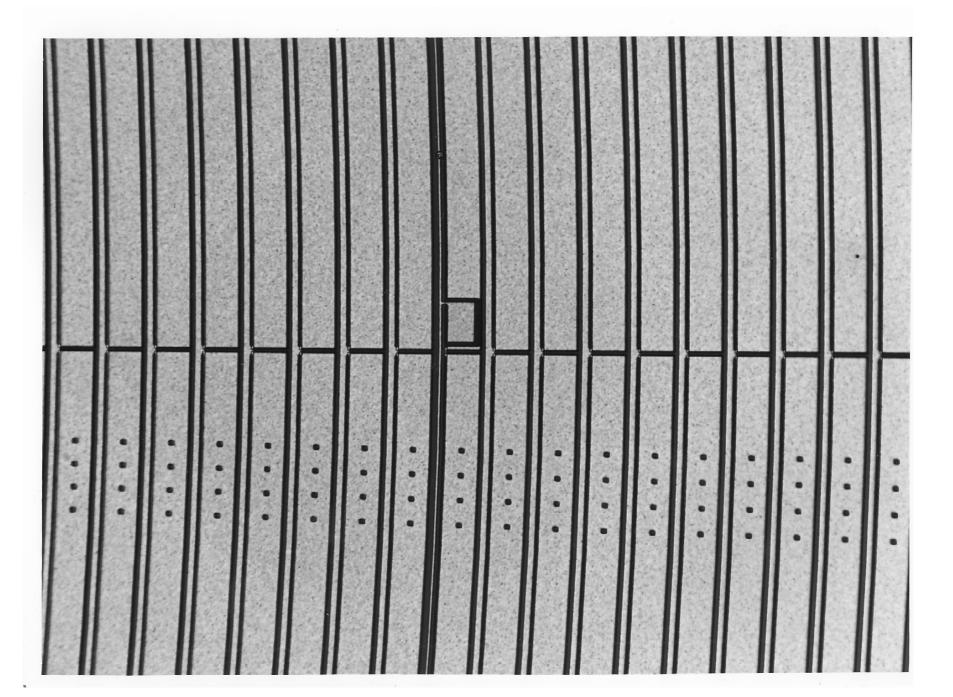
Where: w(t) is unknown weighting function, f(t) is pulse shape induced on the anode normalized to the area of 1, N is the number of electrons in the signal pulse, C total input capacitance,  $e_n^2$  physical spectral density of the series noise, vq is the leakage current and q is the positive value of charge of an electron.

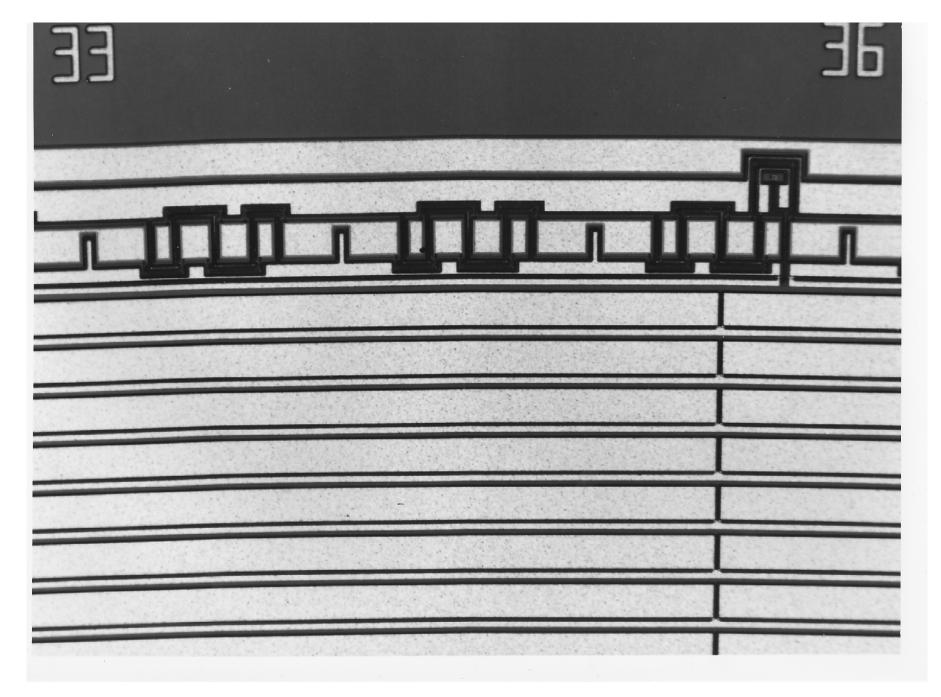
To minimize the above time variance can be reduced to the solution of the following equation:

$$-\frac{C^{2}e_{n}^{2}}{q^{2}N^{2}}w''(t) + \frac{2f(t)}{N}w(t) + \frac{2v}{N^{2}}w(t) = 2\varepsilon_{\min}^{2}f'(t) \cdot const$$

# Position sensing in high energy heavy ion experiments

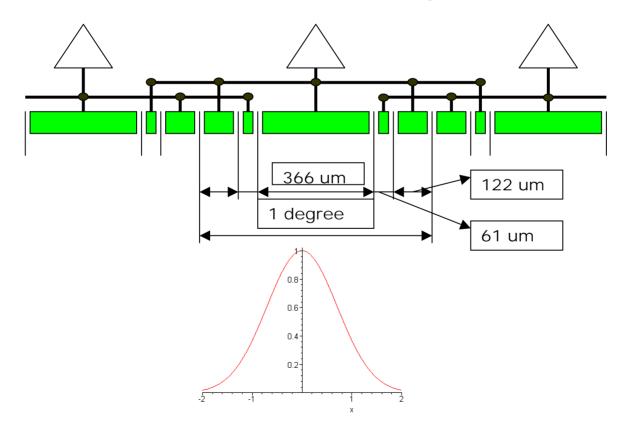
- NA45 (CERES) Experiment at SPS at CERN a) 3" cylindrical drift detector and b) 4" cylindrical detector (past)
- STAR Experiment at RHIC at BNL (present)
- ALICE Experiment at LHC at CERN (future)



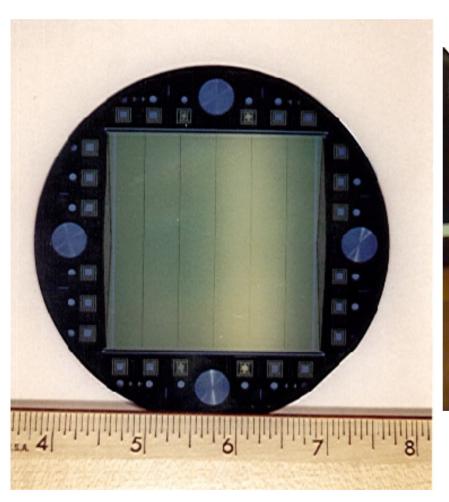


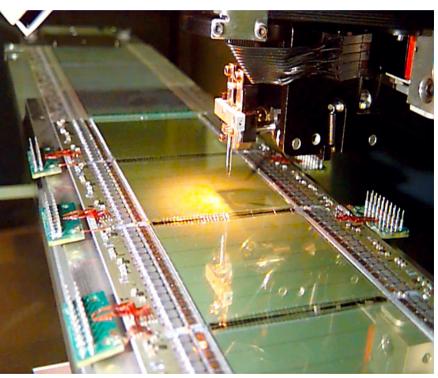
# Innovations in CERES cyl. Det.

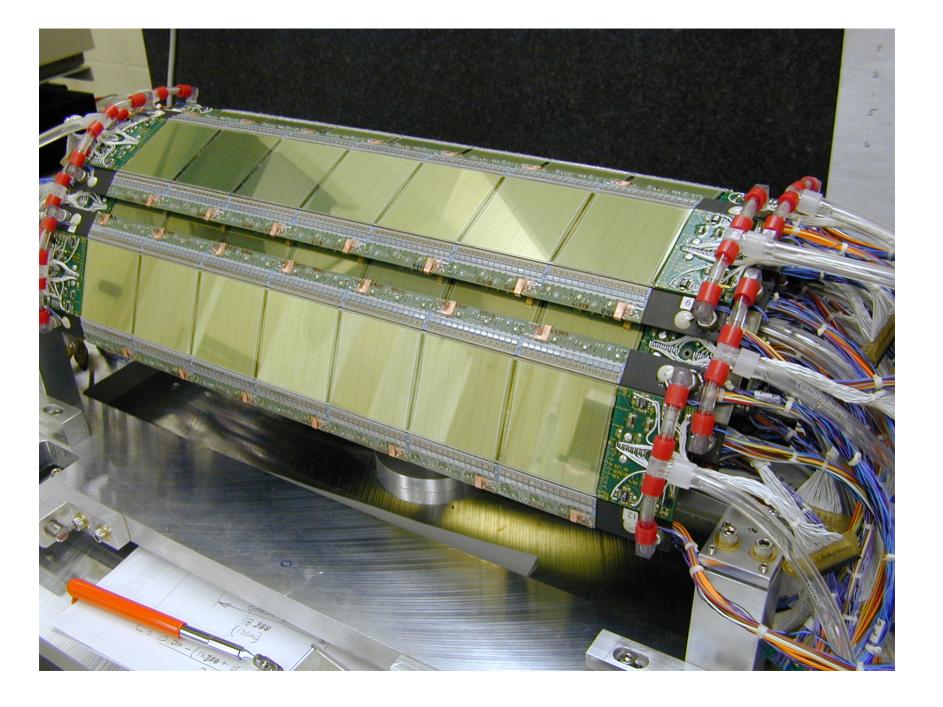
- Collection of leakage current generated at the Si-SiO interface at a sink anode
- Interlaced anodes (Nyquist filtering in linear dimension)



### STAR Drift Detector



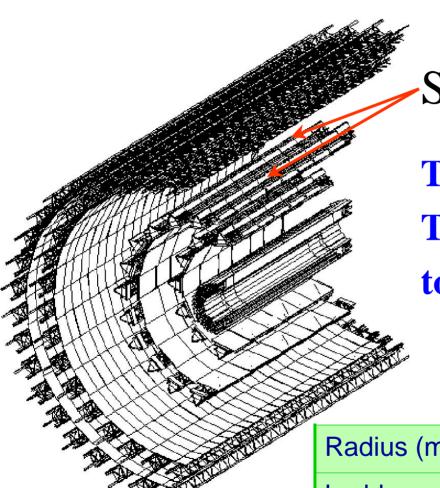




#### SDD collaboration

- INFN Turin Italy
- INFN Trieste Italy
- INFN Bologna Italy
- INFN Rome Italy
- INFN Alessandria- Italy
- Ohio State University Columbus Ohio USA
- University of Jyvaskyla Jyvaskyla Finland
- Nat. Acad. of Sciences, Bogolyubov Inst. for Th. Phys. Kiev -Ukraine
- Scientific Res. Techn. Inst. of Instrument Making Kharkov -Ukraine
- Acad. of Sciences of Czech Republic Rez U Prahy Czech Republic
- St. Petersburg State University St. Petersburg Russia

#### SDD barrels



-Silicon Drift Detectors

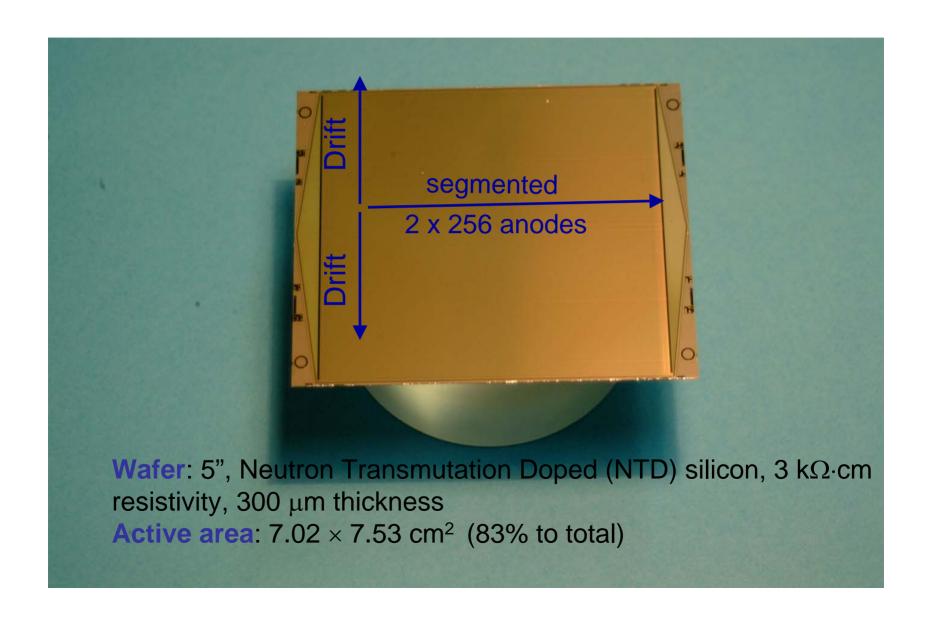
Tot. No. channels  $133 \cdot 10^3$ 

Tot. No. detectors 260

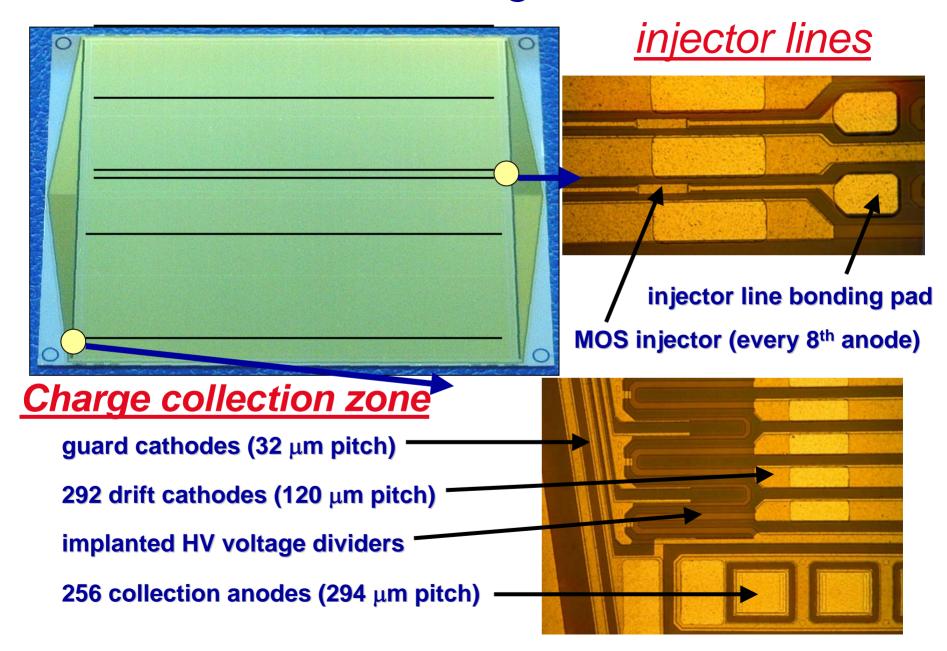
total area 1.37 m<sup>2</sup>

	Layer 3	Layer 4
Radius (mm)	14.9	23.8
Ladders	14	22
SDDs per ladder	6	8

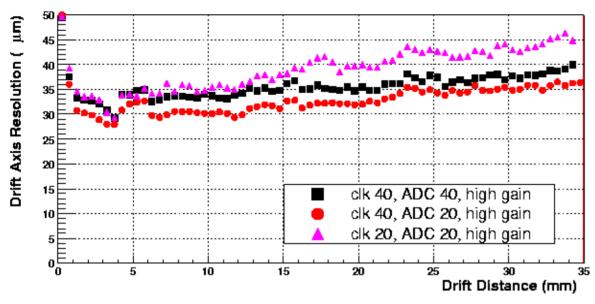
#### ALICE Silicon Drift Detector



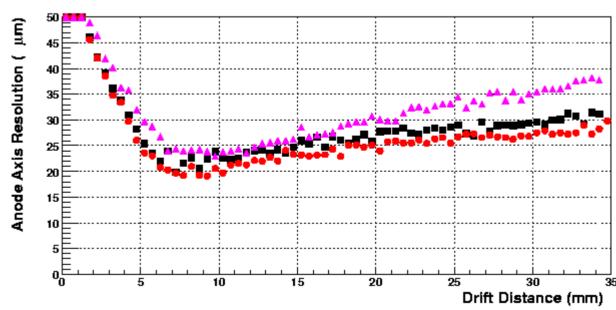
### Detector design features



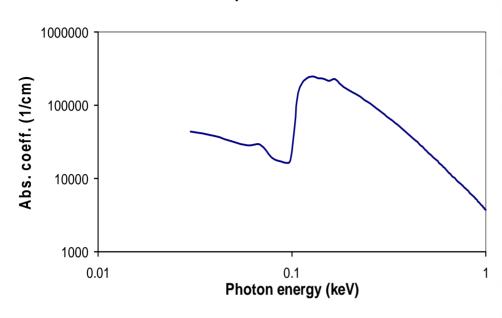
#### Beam Tests in 2003 – Position resolution

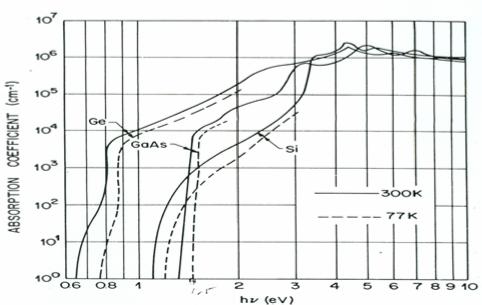


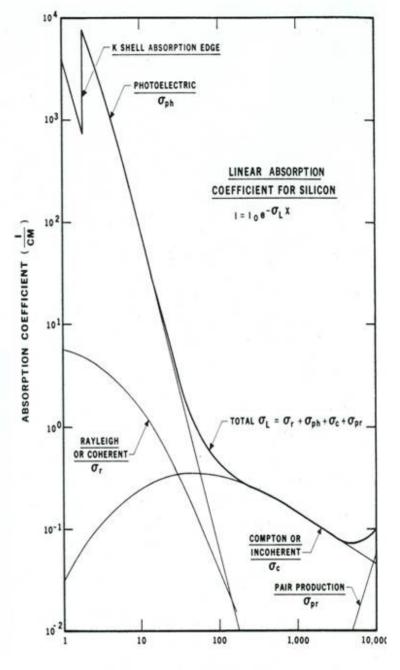
Exhaustive test of the front-end parameters performed in beam test:gain, sampling & ADC frequencies...



#### **Linear absorption coefficient**







$$\sigma_{Si}^2 = F \cdot E_{photon} w$$

$$ENC_{par}^{2} = qI_{leak} \int h(t)^{2} dt \approx qI_{leak} t_{peak}$$

$$ENC_{series}^{2} = \frac{1}{2} e_n^2 C_t^2 \int_{-\infty}^{\infty} [h'(t)]^2 dt \approx e_n^2 C_t^2 / t_{peak}$$

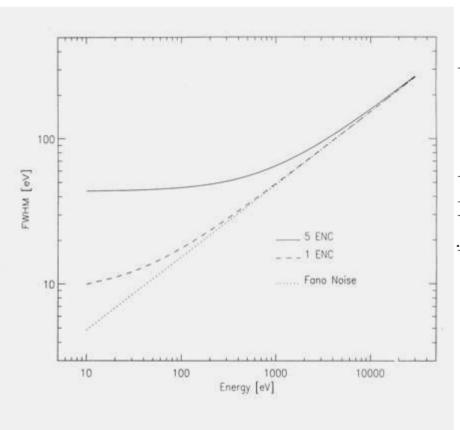
$$t_{peak}^{(optimal)} = e_n C_t / \sqrt{q \cdot I_{leak}}$$

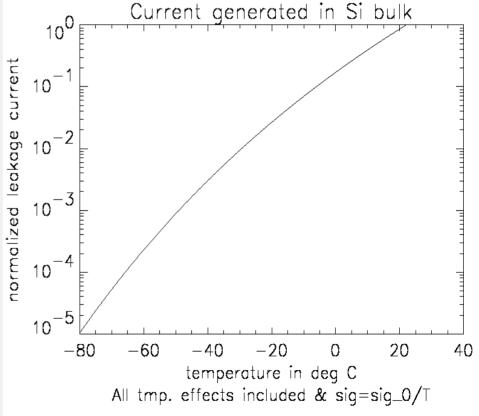
$$w = 3.62eV$$

Leakage Current: 
$$i_{leak} = q \cdot n_i / (2 \cdot \tau)$$

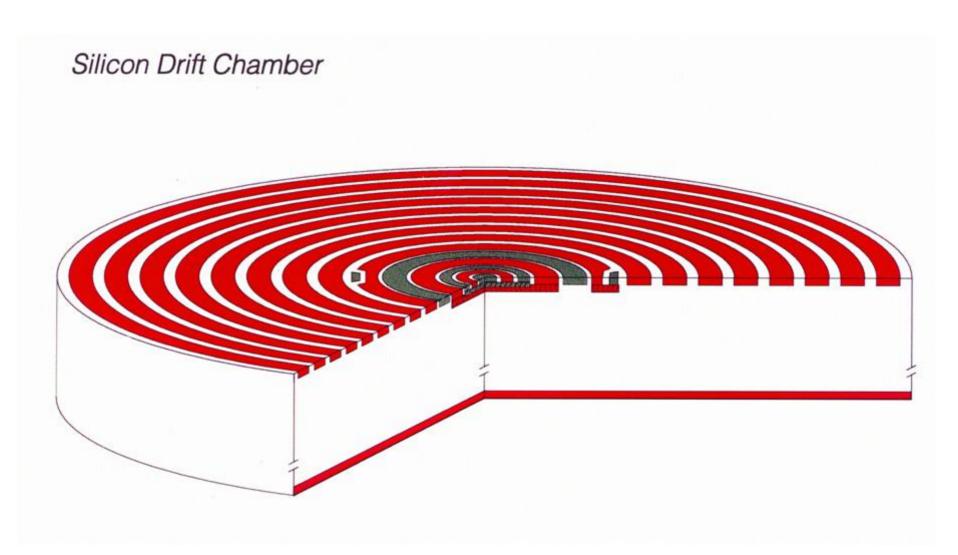
Where n\_i is the density of carrier in intrinsic silicon, tau is the life time and N\_t the density of traps in silicon bulk.

$$n_i \propto \exp(-E_{gap}/(2kT))$$
  
 $\tau \approx 1/(\sigma \cdot v_{th} \cdot N_t)$ 

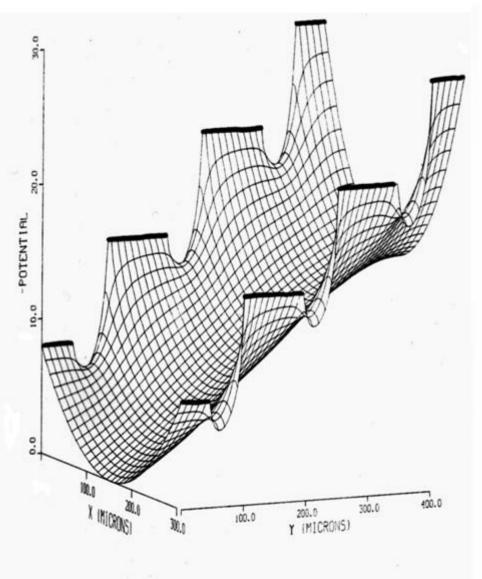


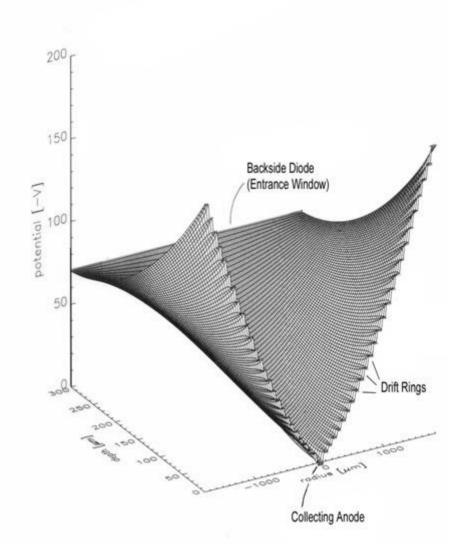


# X-ray drift detector



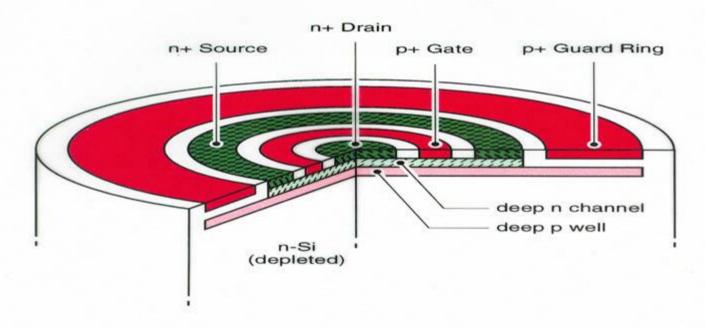
### Potential within Drift Detectors





## Single sided junction FET

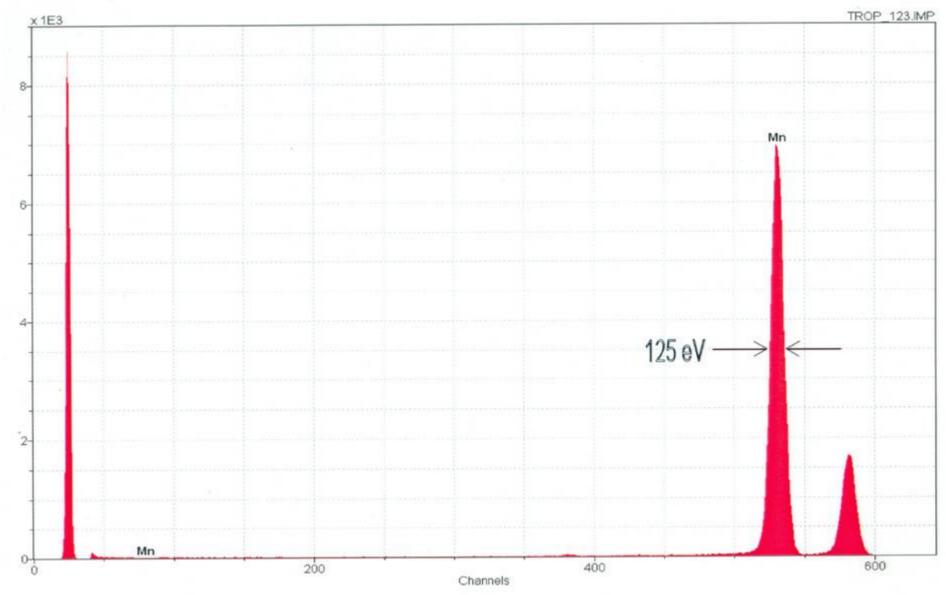
N-Channel JFET on Depleted N-type Silicon



size & characteristics (typical):

gate length	5	μm
gate width	50	μm
saturation current	400	μА
transconductance	400	μA/V

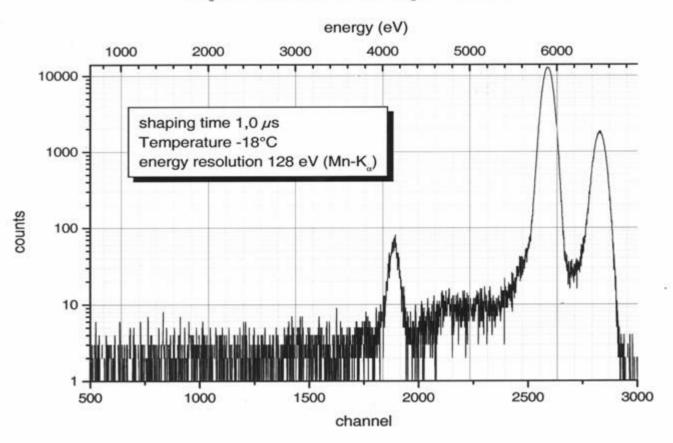
### The best room temp. spectrum



# Low energy tails



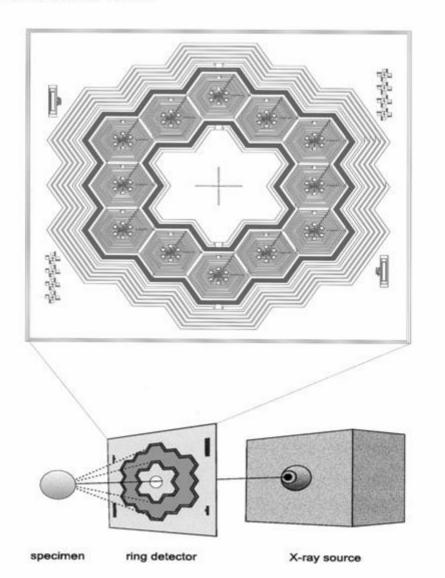
#### Spectral Response



### X-ray fluorescence system

Multi Channel SDD

ring detector for compact XRF systems

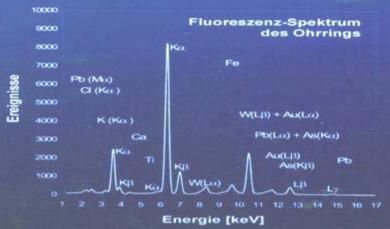


### Applications in Art studies

XRF-Analyse (X-Ray Fluorescence)

Untersuchung eines Leichentuchs (Antinopolis, III. Jahrhundert n.Chr., Vatikanische Museen)



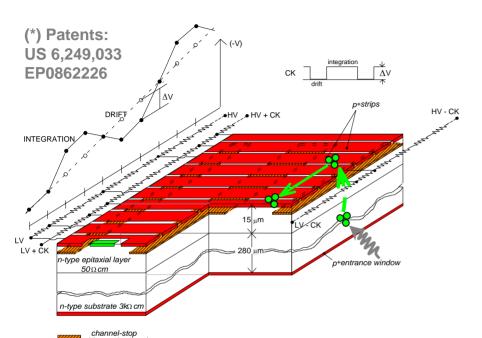




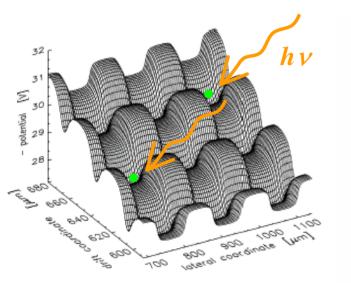
Photographie des Detektor-Moduls

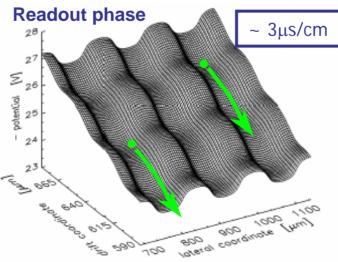
Die Farbe besteht aus einer Mischung von Orpiment (As,S,) und Goldstaub.

#### The Controlled-Drift Detector (CDD)\*



- 2D position sensing (100-200 $\mu$ m)
- low capacitance (~100fF) and integrated JFET⇒ high energy resolution
- low no. of channels (n instead of  $n \times n$ )
- integrate-readout mode





The X-ray position along the drift is obtained from the electrons' drift time

The X-ray energy is obtained from the electron charge collected at the anodes

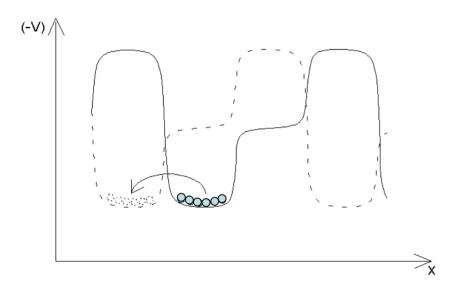
#### Transport mechanism and Readout speed

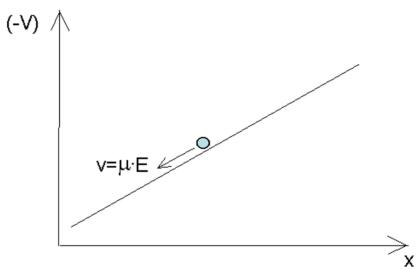
#### **Charge-Coupled Device (CCD)**

L.Strüder et al., NIM A257 (1987) 594

#### **Controlled Drift Detector**

A.Castoldi et al., IEEE TNS, 44(5) oct. 1997 p.1724





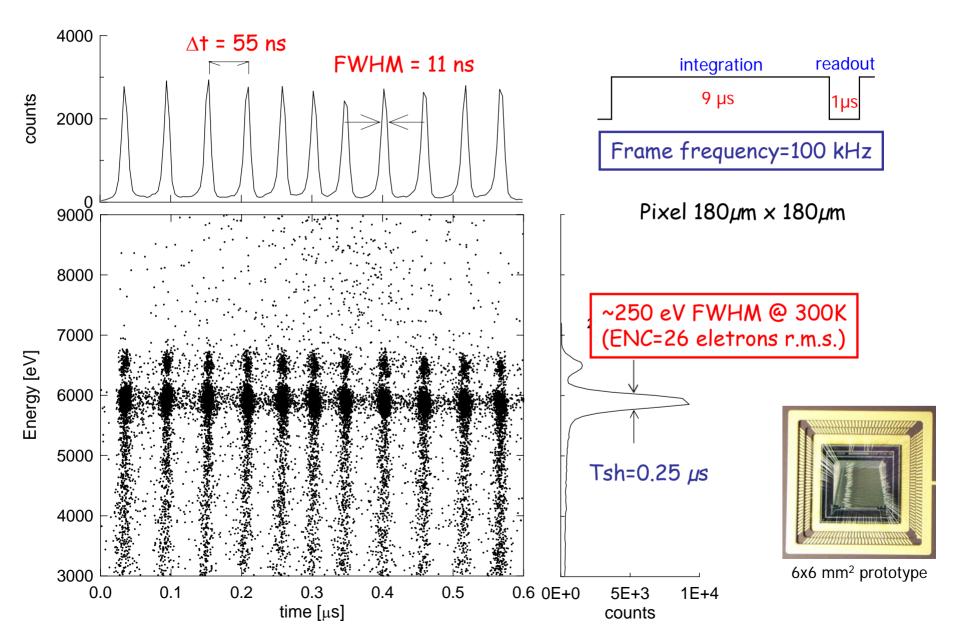
- long readout time as charge transfer and processing are done sequentially
- Shorter readout time as charge transfer and processing are simultaneous

$$T_{readout} = N_{pixel} (T_{proc} + \Delta t_{tr}) \sim 1 \text{ ms/cm}$$

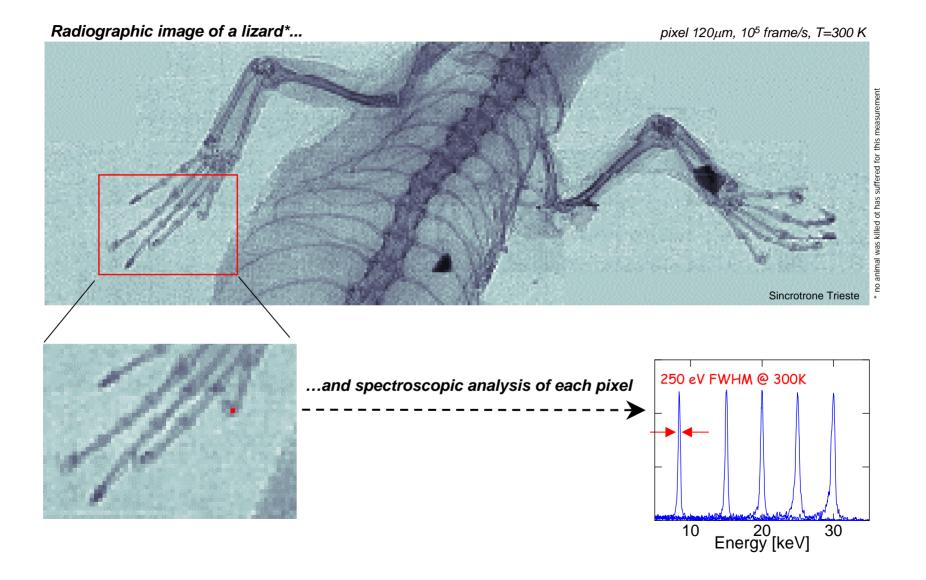
$$T_{readout} = T_{drift}$$
 ~ 3 µs/cm

#### 1-D imaging and spectroscopy of a Fe-55 source @ 100 kHz

A.Castoldi, C.Guazzoni, P.Rehak, L.Strüder, Trans. Nucl. Sci. 49 (3) June 2002



#### X-ray spectroscopic imaging with CDDs

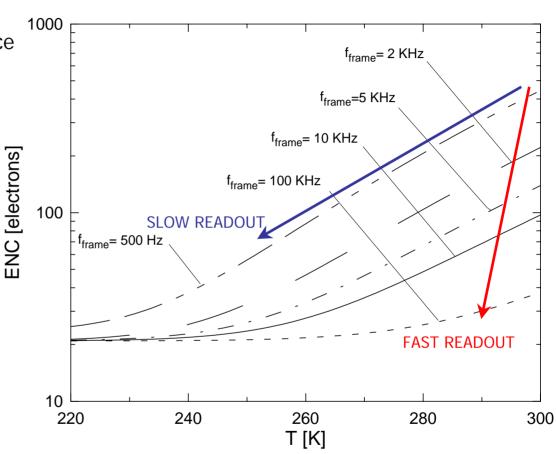


#### Readout speed and Energy resolution

A fast readout speed allows to reduce both readout and integration times:



- higher frame rate (i.e. better time resolution between X-ray images)
- better energy resolution at room temperature due to lower integrated leakage charge.



Time-resolved imaging at frame frequency greater than 10 kHz State-of-the-art energy resolution near room T

## Summary and Conclusions

- Extended use of SDD for tracking in high energy heavy ions experiments
- Industrial use of Silicon Drift Detectors for X-ray fluoroscopy
- Development of Controlled Drift Detector
- Use of Drift concept for Detectors on High Z materials. (insensitive to hole trapping)
- Future-high resistivity silicon for X-rays
- Future tracking with electronics grade Si