## Unfolding Methods

Alberto Rotondi Università di Pavia

# Folding is a common process in physics



### **Convolution is a linear folding**



$$g(y) = \int f(x) \,\delta(y - x) \,\mathrm{d}x \;,$$



### Folding theorem

 $Z = f(X_1, X_2) .$   $Z_1 \equiv Z \ \mathbf{e} \ Z_2 = X_2 \quad \text{auxiliary variable}$  $Z_1 = f(X_1, X_2) , \qquad Z_2 = X_2 .$ 

The jacobian is

$$|J| = \begin{vmatrix} \frac{\partial f_1^{-1}}{\partial z_1} & \frac{\partial f_1^{-1}}{\partial z_2} \\ 0 & 1 \end{vmatrix} = \begin{vmatrix} \frac{\partial f_1^{-1}}{\partial z_1} \\ \frac{\partial f_1^{-1}}{\partial z_1} \end{vmatrix},$$

From the general theorem one obtains

$$p_{\boldsymbol{Z}}(z_1, z_2) = p_{\boldsymbol{X}}(x_1, x_2) \left| \frac{\partial f_1^{-1}}{\partial z_1} \right|$$

by integrating on the auxiliary variable :  $p_{Z_1}(z_1) = \int p_{\boldsymbol{Z}}(z_1, z_2) \, \mathrm{d} z_2$  .

hence

$$p_Z(z) = \int p_{\boldsymbol{X}}(x_1, x_2) \left| \frac{\partial f_1^{-1}}{\partial z} \right| \, \mathrm{d}x_2$$
$$= \int p_{\boldsymbol{X}} \left( f_1^{-1}(z, x_2), x_2 \right) \left| \frac{\partial f_1^{-1}}{\partial z} \right| \, \mathrm{d}x_2 \,,$$

which is the probability density

For independent variables:

$$p_Z(z) = \int p_{X_1} \left( f_1^{-1}(z, x_2) \right) \, p_{X_2}(x_2) \frac{\partial f_1^{-1}}{\partial z} \, \mathrm{d}x_2 \, .$$

## Convolution theorem

when Z is given by the sum

$$Z = X_1 + X_2 ,$$

#### we have

$$X_1 = f_1^{-1}(Z, X_2) = Z - X_2 , \qquad \frac{\partial f_1^{-1}}{\partial z} = 1 ,$$

and we obtain

$$p_Z(z) = \int_{-\infty}^{+\infty} p_X(z - x_2, x_2) \,\mathrm{d}x_2$$
.

When  $X_1$  and  $X_2$  are independent, we obtain the convolution integral

$$p_Z(z) = \int_{-\infty}^{+\infty} p_{X_1}(z - x_2) \, p_{X_2}(x_2) \, \mathrm{d}x_2 \; ,$$

In physics

(  $\delta$  instrument function, f signal)  $g(y) = \int_{-\infty}^{+\infty} f(y - x) \,\delta(x) \, \mathrm{d}x ,$ 

### **Uniform\*Gaussian**

When Z = X + Y where  $X \sim N(\mu, \sigma^2)$  e  $Y \sim U(a, b)$ . one has immediately

$$p_Z(z) = \frac{1}{b-a} \int_a^b \frac{1}{\sigma\sqrt{2\pi}} \exp\left[-\frac{(z-y-\mu)^2}{2\sigma^2}\right] dy = \frac{1}{b-a} \int_a^b \frac{1}{\sigma\sqrt{2\pi}} \exp\left[-\frac{[y-(z-\mu)]^2}{2\sigma^2}\right] dy .$$

$$p_Z(z) = \frac{1}{b-a} \left[ \Phi\left(\frac{b-(z-\mu)}{\sigma}\right) - \Phi\left(\frac{a-(z-\mu)}{\sigma}\right) \right]$$





In the reconstruction of an histogram,

• the true histogram (image) where the bin contents are the expected values

$$oldsymbol{\mu} = (\mu_1, \mu_2, \dots, \mu_N) \;, \;\; \mu_j = \mu_{ ext{tot}} p_j = \mu_{ ext{tot}} \int_{ ext{bin i}} f_t(y) \, \mathrm{d} y$$



The observed  $N_{ij}(exp)$  events have to be compared with the expected values  $N_{ij}(th)$  predicted by a model.

$$N_{ij}(\text{th}) = NP_{ij}(\text{obs}) = N\sum_{i'j'} P_{i'j'}(\text{true}) P_v(\text{obs}_{ij}|\text{true}_{i'j'}) , \quad (94)$$

that is, the number of events observed in the ijth-cell is due to the presence into the i'j'th-cell, times the probability  $P_v$  that the PSF shifts the point from the i'j' to the ij-cell. One has to sum on all the cells near the ij-one.

In eq. (94) the normalization is understood. In practice, from (93, 94), one has

$$P(\operatorname{true}_{ij}|\operatorname{obs}_{i'j'}) = \frac{P_{i'j'}(\operatorname{true}) P_v(\operatorname{obs}_{ij}|\operatorname{true}_{i'j'})}{\sum_{ij} \left[\sum_{i',j'} P_{i'j'}(\operatorname{true}) P_v(\operatorname{obs}_{ij}|\operatorname{true}_{i'j'})\right]} \quad (95)$$

In the case of a two dimensional Gaussian point spread function **PSF**:

$$P_v(\text{obs}_{ij}|\text{true}_{i'j'}) = \frac{1}{2\pi \,\sigma_x \sigma_y} \,\exp\left[-\frac{(x_{ij} - x_{i'j'})^2}{2 \,\sigma_x^2} - \frac{(y_{ij} - y_{i'j'})^2}{2 \,\sigma_y^2}\right],$$
(96)

### Fourier Techniques

$$f(x) = \int F(t) e^{2\pi i xt} dt$$

**Convolution:** 

$$f(x) = \int g(y)\delta(x-y) \,\mathrm{d}y$$
$$\int F(t) \,\mathrm{e}^{2\pi i xt} \,\mathrm{d}t = \int G(t) \,\,\mathrm{e}^{2\pi i yt} \,\,\Delta(t) \,\mathrm{e}^{2\pi i (x-y)t} \,\,\mathrm{d}t$$
$$\int F(t) \,\mathrm{e}^{2\pi i xt} \,\,\mathrm{d}t = \int G(t) \,\,\Delta(t) \,\mathrm{e}^{2\pi i xt} \,\,\mathrm{d}t \to F(t) = G(t) \,\,\Delta(t)$$

### **Image Deconvolution**

$$D(\boldsymbol{x}) = \int \mathrm{d}\boldsymbol{y} I(\boldsymbol{y}) \,\delta(|\boldsymbol{x} - \boldsymbol{y}|)$$

In the absence of noise

$$I = F^{-1} \left[ \frac{F(D)}{F(\delta)} \right]$$

where F is the Fourier transform.

For a real image  $I(n_1, n_2)$  the Fourier transform is:

$$F(k_1, k_2) = \sum_{n_2=0}^{N_2-1} \sum_{n_1=0}^{N_1-1} e^{2\pi i k_2 n_2/N_2} e^{2\pi i k_1 n_1/N_1} I(n_1, n_2)$$

$$F(k_1, k_2) = FFT_2[FFT_1[I(n_1, n_2)]]$$
or the routines see for example Numeric

For the routines see for example *Numerical Recipes* 

### The problem with fluctuations

Inverting the response mitrix 161



Fig. 11.1 (a) A hypothetical true histogram  $\mu$ . (b) the histogram of expectation values  $\nu = R\mu$ . (c) the histogram of observed data n. and (d) the estimators  $\hat{\mu}$  obtained from inversion of the response matrix.



#### Poisson statistics

Fourier restored

Figure 11: Lena restored by FFT: The original image (top left) is sampled with Poisson statistics (top right) and smeared with a 2D 10-bins Gaussian PSF (bottom left): the Fourier restored image (bottom right) is similar to the Poisson sampled image. In this case the noise term N is neglected.

#### original

#### Poisson statistics



#### Fourier (un)restored

Gaussian

smearing

Figure 12: Lena not restored by FFT: In this case the noise term N is not ignored: the original image (top left) is smeared with a 2D 10-bins Gaussian PSF (top right) and the result is sampled with Poisson statistics (bottom left): the Fourier restored image (bottom right) cannot recover the information lost in the noise. Another approach, statistical in nature, is required.



Gaussian

smearing

#### einstein einstein 1 1 1 1 7 1 . . . . . . . . . . . . . . . einstein .0 Ð

Poisson statistics

Fourier restored

Figure 13: Einstein restored by FFT: explanation as in Figure 1.



Figure 14: Einstein not restored by FFT: explanation as in Figure 2.

#### Fourier (un)restored

### PET: positron emission thomography















When true  $\rightarrow$  obs  $\equiv \mu \rightarrow \nu$  deterministic methods (FFT) can be used

$$\boldsymbol{\nu} = R \ast \boldsymbol{\mu} \tag{19}$$

When true  $\rightarrow$  smeared  $\rightarrow$  obs  $\equiv \mu \rightarrow \nu \rightarrow n$  statistical methods must be used

$$n = \nu + \rho = R * \mu + \rho$$

In the poissonian or binomial case we have to minimize:

$$-\ln L(\boldsymbol{\mu}) = -\sum_{i} \ln P(n_i, \nu_i)$$

In the gaussian case we must minimize

$$\chi^{2}(\boldsymbol{\mu}) = \sum_{ij} (\nu_{i} - n_{i}) (V^{-1})_{ij} (\nu_{j} - n_{j})$$

In 2-D, when  $N_{ij}(exp)$  contains fluctuations, we have to minimize:

$$-2 \ln L(\boldsymbol{n}|\boldsymbol{\nu}, \boldsymbol{\mu}) \simeq \chi^2 = \sum_{ij} \frac{[N_{ij}(\exp) - NP_{ij}(\operatorname{obs})]^2}{NP_{ij}(\operatorname{obs})}$$
(20)

where

$$P(\text{obs}) = P(\nu|\mu) P(\mu)$$

If all the pixel contents  $\mu$  are the free parameters to be determined the problem has zero DoF This is a ILL-POSED problem with many (and more probable) unrealistic solution!!

## Image restoration

Unlike!

### **Explanation:**



many solutions give a good  $\chi^2$ the spike ones are more probable! Cure: to add to  $\chi^2$  an empirical regularization term C[p].

 $\chi^2 \to \alpha \, \chi^2 + C[P(\text{true})]$ 

 $\mathbf{or}$ 

 $\chi^2 \to \chi^2 + \alpha C[P(\text{true})]$ 

or

to increase the DoF by using a parametric model  $P(\nu \mid \mu)P(\mu) \rightarrow P(\nu \mid \mu')$  Remember

$$\mu_i \stackrel{PSF}{\rightarrow} \nu_i \stackrel{random}{\rightarrow} n_i$$

and consider (95) as a form of the Bayes theorem

$$P(\text{true}_{ij}|\text{obs}) \propto \sum_{i',j'} P_{i'j'}(\text{true}) P_v(\text{obs}_{ij}|\text{true}_{i'j'}) = L(\boldsymbol{n}|\boldsymbol{\mu})P(\boldsymbol{\mu})$$

Bayesians say: posterior = likelihood × prior One maximizes  $P(\text{true}_{ij}|\text{obs}) \equiv F(\boldsymbol{\mu})$  (or minimize  $-F(\boldsymbol{\mu})$ ):

$$F(\boldsymbol{\mu}) = \ln L(\boldsymbol{n}|\boldsymbol{\mu}) + \ln P(\boldsymbol{\mu})$$
(99)

following the Maximum Likelihood (ML) principle. The practical (no Bayesian) experimentalist introduces an empirical regularization parameter  $\alpha$  and considers the prior  $P(\mu)$  as a regularization function  $C(\mu)$ : The frequentist assumes

$$P(\mu) = 1$$

$$F(\boldsymbol{\mu}) = \alpha \ln L(\boldsymbol{n}|\boldsymbol{\mu}) + C(\boldsymbol{\mu})$$
(100)

By keeping fixed the normalization:

$$\nu_T = \sum_i \sum_j R_{ij}\hat{\mu}_j + \rho_i = n_T$$

the objective function is

$$F(\boldsymbol{\mu}) = \alpha \ln L(\boldsymbol{n}|\boldsymbol{\mu}) + C(\boldsymbol{\mu}) + \lambda(n_T - \sum_i \nu_i) \quad (101)$$

where  $\lambda$  is a Lagrange multiplier

$$\frac{\partial F}{\partial \lambda} = 0 \rightarrow \sum_{i} n_{i} = n_{T}$$

The objective function to be minimized is

$$-F(\boldsymbol{\mu}) = -2\ln L(\boldsymbol{n}|\boldsymbol{\mu}) - \alpha \, \boldsymbol{C}(\boldsymbol{\mu}) + \lambda(n_T - \sum_i \nu_i) \quad (23)$$

$$\mu_j = \mu_{\text{tot}} p_j = \mu_{\text{tot}} \int_{\text{bin j}} f_t(y) \, \mathrm{d}y$$

where  $\alpha > 0$ . Some regularization terms:

• minimum second derivative (Tichonov)

$$C(\boldsymbol{\mu}) = -\int [f_t''(y)]^2 \, \mathrm{d}y \simeq -\sum_{i=1}^{M-2} [-\mu_i + 2\mu_{i+1} - \mu_{i+2}]^2$$

• minimum variance:

$$C(\boldsymbol{\mu}) = -\operatorname{Var}[\boldsymbol{\mu}] \equiv ||C\boldsymbol{\mu}||^2 = -\sum_i \mu_i^2$$

• maximum entropy (MaxEnt)

$$C(\boldsymbol{\mu}) = -\sum_{i} p_{i} \ln p_{i} = -\sum_{i} \frac{\mu_{i}}{\mu_{T}} \ln \frac{\mu_{i}}{\mu_{T}}$$

cross-entropy

$$C(\boldsymbol{\mu}) = -\sum_{i} p_i \ln \frac{p_i}{q_i} = -\sum_{i} \frac{\mu_i}{\mu_T} \ln \frac{\mu_i}{\mu_T q_i}$$

where  $q = (q_1, q_2, \ldots, q_n)$  is the most likely a priori shape for the true distribution  $\mu_i$ .

## Regularization terms

We have M boxes and a monkey that throws N ball randomly into them.

What is the box-balls configuration of highest probability? Probability of a configuration:



### What is MaxEnt ???

23

$$N_{ij}(\text{th}) = NP_{ij}(\text{obs}) = N \sum_{i'j'} P_{i'j'}(\text{true}) P_v(\text{obs}_{ij}|\text{true}_{i'j'}) \quad (25)$$
  
We write this equation considering the operator  $R$ :  
$$n - R * \mu$$
  
The iterative method (Van Cittert 1930) add, with a weight  $i \rightarrow 0$  residual  $r_k$  to the current solution  
$$\mu_{k+1} = \mu_k + \beta[n - R * \mu_k] \quad (26)$$
  
The method is based on the known equation  
$$\sum_{i=0}^k q^n = \frac{1 - q^{k+1}}{1 - q} \quad (27)$$

for  $k \to \infty$  the series converges if |q < 1|. By applying iteratively (26)

$$\mu_{k+1} = \beta n + (1 - \beta R) \mu_k = \beta n + (1 - \beta R) (\beta n + (1 - \beta R) \mu_{k-1})$$
  
=  $\beta n + \beta (1 - \beta R) n + (1 - \beta R)^2 \mu_{k-1}$   
=  $\beta n + \beta (1 - \beta R) n + \beta (1 - \beta R)^2 n + (1 - \beta R)^3 \mu_{k-2} \dots$   
=  $\sum_{i=0}^k \beta (1 - \beta R)^i n$ .

From (27):

$$\mu_{k+1} = \frac{I - (I - \beta R)^{k+1}}{\beta R} \beta n \to R^{-1} n = \mu , \quad \text{for} \quad k \to \infty .$$

if  $|I - \beta R| < 1$ 

Without 
$$\beta$$

 $\mu_{k+1} = \mu_k + [n - R * \mu_k]$ 







### The iterative Principle without best fit + smoothing



Consider the case with statistical fluctuations

$$\mu_i \stackrel{PSF}{\to} \nu_i \stackrel{random}{\to} n_i$$
$$\mathbf{n} = \mathbf{R} * \mu + \mathbf{r}$$

To have regular solutions and to make the method robust, we must search for an iterative solution which minimize the  $\chi^2$ :

$$\chi^2 = ||\mathbf{R} * \mu - \mathbf{n}||^2 = \frac{1}{2} \sum_{ik} (\sum_{mn} \mathbf{R}_{i-m,k-n} \mu_{mn} - \mathbf{n}_{ik})^2$$

Note that is the case in which

the PSF depends on the pixel difference only (translational invariance)

In this case we consider R as an operator and we can work with symmetric  $M \times N$  matrices.

Minimum  $\chi^2$  w.r.t  $\mu_{ik}$  gives the equations:



The iterative algorithm + best fit It belongs to the family of the Newton-Raphson methods From (29):



This formula minimizes  $\chi^2$ . For example when  $\beta_k = \beta$ :

$$\begin{aligned} \mu_{k+1} &= \mu_k + \beta R * [n - R * \mu_k] \\ &= \beta R * n + (I - \beta R^2) * \mu_k \\ &= \beta R * n + \beta (I - \beta R^2) * R * n + (I - \beta R^2)^2 * \mu_{k-1} \\ &= \sum_{i=0}^k \beta (I - \beta R^2)^i * R * n = \frac{I - (I - \beta R^2)^{k+1}}{\beta R^2} \beta R * n \to R^{-1}n \end{aligned}$$

The iterative algorithm + Best fit

29

$$\mu_{k+1} = \mu_k + \beta_k R * [n - R * \mu_k]$$

Previous method converges if

$$||I - \beta R * R|| < 1$$

when  $\beta$  is independent of k. In this case

$$0 < \beta < \frac{2}{\text{max eigenvalue of } (R * R * \mu) * \mu^{-1}}$$

The iterative algorithm + best fit + regularization

(31)

When  $\beta_k$  depends on k convergence is assured if (Robbins and Munro 1951)

$$\lim_{N \to \infty} \beta_N = 0 , \quad \sum_{N=1}^{\infty} \beta_N = \infty , \quad \sum_{N=1}^{\infty} \beta_N^2 < \infty ,$$

Next, the method must be regularized by adding a term to the  $\chi^2$ .  $\chi^2 = ||R * \mu - n||^2 + \alpha ||C * \mu||^2$ (32) For example,  $||\mathcal{C} * \mu||^2 = \sum_{ik} \mu_{ik}^2$  $||C * \mu||^2 = |\sum_{ik} \mu_{ik} \ln \mu_{ik} / \mu_{tot}| , \quad \alpha < 0$ 

The iterative solution becomes

$$\mu_{k+1} = \mu_k + \beta_k [R * n - (R * R + \alpha C * C)\mu_k]$$
(33)

$$\mu_{k+1} = \mu_k + \beta_k [R * n - [R * R * \mu_k + \alpha(\ln \mu_k / \mu_T + I)]$$

About 40 iterations, regularized with Maximum entropy



The iterative algorithm + best fit + MaxEnt regularization

$$\mu_{k+1} = \mu_k + \beta_k [R * n - (R * R + \alpha I) * \mu_k]$$

About 100 iterations, regularized with the sum of squares



The iterative algorithm + best fit + Tichonov regularization

$$\mu_{k+1} = \mu_k + \beta_k [R * n - (R * R + \alpha I) * \mu_k]$$



The iterative algorithm + best fit + Tichonov regularization

$$\mu_{k+1} = \mu_k + \beta_k [R * n - (R * R + \alpha I) * \mu_k]$$



The iterative algorithm + best fit + Tichonov regularization

### **ATHENA** apparatus



### From the ATHENA detector





Distribution of annihilation vertices when antiprotons are mixed with ...

Pbar-only (with electrons)



### hot positrons





antihydrogen !!!!!!!!!

#### **FIRST COLD ANTIHYDROGEN PRODUCTION & DETECTION (2002)** M. Amoretti et al., Nature 419 (2002) 456 M. Amoretti et al., Phys. Lett. B 578 (2004) 23



#### Annihilation vertex in the trap x-y plane



BCKG

**Hbar percentage**  $x = 0.65 \pm 0.05$ 



### Iteratve best fit method



Cold Mix

The 2D deconvolution reveals two different annihilation modes The iterative algorithms + best fit + regularization

- iterative algorithms are used in unfolding (ill posed) problems
- they need a Bayesian regularization term
- when there are degrees of freedom, one can use a best fit of a signal+background function to the data
- in this case there are no Bayesian terms (pure frequentist approach)



### Conclusions

- don't be dogmatic
- use Bayes to parametrize the a priori knowledge if any, not the ignorance
- in the case of poor a priori knowledge, use the frequentist methods

The objective function to be minimized is

$$-F(\boldsymbol{\mu}) = -2\ln L(\boldsymbol{n}|\boldsymbol{\mu}) - \boldsymbol{\alpha} C(\boldsymbol{\mu}) + \lambda(n_T - \sum_i \nu_i) \quad (24)$$

Some choices of  $\alpha > 0$  are:

• Bayesian

$$\alpha = \frac{1}{\mu_T}$$

usually too much smoothing

•  $A = \chi^2 + \alpha ||C\mu||^2 > 0$ 

we can regularize the solution by choosing  $\chi^2 \simeq DoF \equiv$  number of pixel N, with the condition:

$$\alpha = \frac{A - N}{||C\mu||^2}$$



## Regularization parameter

#### The problem with fluctuations

From  $\nu \rightarrow \mu$  deterministic methods can be used

$$\boldsymbol{\nu} = R \,\boldsymbol{\mu} \tag{97}$$

From  $n \rightarrow \nu \rightarrow \mu$  statistical methods must be used

$$\boldsymbol{n} = \boldsymbol{\nu} + \boldsymbol{\rho} = R \, \boldsymbol{\mu} + \boldsymbol{\rho}$$

In the poissonian or binomial case we have to minimize:

$$-\ln L(\boldsymbol{\mu}) = -\sum_{i} \ln P(n_i, \nu_i)$$

In the gaussian case we must minimize

$$\chi^{2}(\boldsymbol{\mu}) = \sum_{ij} (\nu_{i} - n_{i})(V^{-1})_{ij}(\nu_{j} - n_{j})$$

These estimators are unbiased:

$$E[\hat{\mu}_j] = \sum_i (R^{-1})_{ji} E[n_i - \rho_i] = \sum_i (R^{-1})_{ji} \nu_i = \mu_j$$

In 2-D, when  $N_{ij}(exp)$  contains fluctuations, we have to minimize:

$$\chi^2 = \sum_{ij} \frac{[N_{ij}(\exp) - NP_{ij}(\operatorname{obs})]^2}{NP_{ij}(\operatorname{obs})}$$
(98)

where

$$P(\text{obs}) = P(\nu|\mu) P(\mu)$$

44

#### Fourier techniques

$$f(x) = \int F(t) e^{2\pi i xt} dt$$

**Convolution:** 

$$f(x) = \int g(y)\delta(x-y) \, \mathrm{d}y$$
$$\int F(t) \, \mathrm{e}^{2\pi i x t} \, \mathrm{d}t = \int G(t) \, \mathrm{e}^{2\pi i y t} \, \Delta(t) \, \mathrm{e}^{2\pi i (x-y)t} \, \mathrm{d}t$$
$$\int F(t) \, \mathrm{e}^{2\pi i x t} \, \mathrm{d}t = \int G(t) \, \Delta(t) \, \mathrm{e}^{2\pi i x t} \, \mathrm{d}t \to F(t) = G(t) \, \Delta(t)$$

Correlation

$$\operatorname{Corr}(g,\delta) \equiv \int g(x+y)\,\delta(y)\,\mathrm{d}y \to G(t)\Delta^*(t)$$

if the functions are **real** 

$$G(t) = G(-t)^*$$
,  $\operatorname{Corr}(g, \delta) \to G(t)\Delta(-t)$ 

Autocorrelation (Wiener theorem)

$$\operatorname{Corr}(g,g) \to |G(t)|^2$$

**Total Power:** 

$$P(f) \equiv \int |f(x)|^2 \,\mathrm{d}x = \int |F(t)|^2 \,\mathrm{d}t$$

**Power Spectral Density** (in the Fourier space):

$$PSD(f) \equiv |F(t)|^2 + |F(-t)|^2 \xrightarrow{f(x) \text{ real }} 2|F(t)|^2 \quad 0 \le t \le \infty$$

## Conclusions

- best fit minimization methods are crucial in physics. They are mainly frequentist
- they are based on the ML and LS algorithms (they are implemented in the ROOT-MINUIT framework)
- to judge the quality of the result, frequentists use the  $\chi^2$  test bayesians use the hypothesis probability
- Bayesian a priori hypotheses should be used with informative priors!!!

$$\mu_i \stackrel{PSF}{\to} \nu_i \stackrel{random}{\to} n_i$$
$$L(\boldsymbol{n}|\boldsymbol{\nu}, \boldsymbol{\mu}) \ P(\boldsymbol{\mu})$$

Bayesians say: posterior = likelihood × prior One maximizes  $P(\text{true}_{ij}|\text{obs}) \equiv F(\boldsymbol{\mu})$  (or minimize  $-F(\boldsymbol{\mu})$ ):

$$F(\boldsymbol{\mu}) = \ln L(\boldsymbol{n}|\boldsymbol{\nu},\boldsymbol{\mu}) + \ln P(\boldsymbol{\mu})$$
(21)

following the Maximum Likelihood (ML) principle. The practical (no Bayesian) experimentalist introduces an empirical regularization parameter  $\alpha$  and considers the prior  $P(\mu)$  as a regularization function  $C(\mu)$ :

$$F(\boldsymbol{\mu}) = \ln L(\boldsymbol{n}|\boldsymbol{\nu},\boldsymbol{\mu}) + \alpha C(\boldsymbol{\mu})$$
(22)

The present status of the Bayesian-Frequentist dispute:

- The The Bayesian: always choose o noninformative prior; reject the concept of the ensemble of identical experiments and give a probability to all the hypotheses.
- The (frequentists) physicists: usually I avoid to give any probability to my hypotheses: they are rejected by the experience (falsification). I use only informative priors into the analysis when certain *a priori information* is available.

## Image restoration

### **2D Unfolding**

A picture in a x - y plane is the result of a double dimensional folding, where the true points are smeared out by detector effects.

$$N = \sum_{ij=1}^{n_c} N_{ij}(\exp) , \qquad (93)$$

N is the total number of events and  $N_{ij}(\exp)$  is the recorded number of event in the pixel placed at the *i*th-row and *j*th-column.

The observed  $N_{ij}(exp)$  events have to be compared with the expected values  $N_{ij}(th)$  predicted by a model.

$$N_{ij}(\text{th}) = NP_{ij}(\text{obs}) = N\sum_{i'j'} P_{i'j'}(\text{true}) P_v(\text{obs}_{ij}|\text{true}_{i'j'}) , \quad (94)$$

that is, the number of events observed in the ijth-cell is due to the presence into the i'j'th-cell, times the probability  $P_v$  that the PSF shifts the point from the i'j' to the ij-cell. One has to sum on all the cells near the ij-one.

In eq. (94) the normalization is understood. In practice, from (93, 94), one has

$$P(\operatorname{true}_{ij}|\operatorname{obs}_{i'j'}) = \frac{P_{i'j'}(\operatorname{true}) P_v(\operatorname{obs}_{ij}|\operatorname{true}_{i'j'})}{\sum_{ij} \left[\sum_{i',j'} P_{i'j'}(\operatorname{true}) P_v(\operatorname{obs}_{ij}|\operatorname{true}_{i'j'})\right]} \quad (95)$$

In the case of a two dimensional Gaussian point spread function **PSF**:

$$P_{v}(\text{obs}_{ij}|\text{true}_{i'j'}) = \frac{1}{2\pi \,\sigma_{x}\sigma_{y}} \exp\left[-\frac{(x_{ij} - x_{i'j'})^{2}}{2 \,\sigma_{x}^{2}} - \frac{(y_{ij} - y_{i'j'})^{2}}{2 \,\sigma_{y}^{2}}\right],$$
(96)

48

In summary, we use

 $\mu_{k+1} = \mu_k + \beta[n - R * \mu_k]$ 

with the initial condition

 $\mu_0=n$  .

The convergence is assured if

$$|I - \beta R| < 1$$

(28)

# The iterative principle

Since  $|1 - \beta x| < 1$  implies  $0 < \beta < 2/x$ , in the case of the operator R, which can be transformed in a square matrix

$$R' = (R \ast \mu) \, \mu^{-1}$$

we obtain the condition:

$$0 < \beta < \frac{2}{\text{max eigenvalue of } \mathbf{R}'}$$

Note that we works always with square matrices  $R * \mu$ ,

 $\mu$ ,  $\mu^{-1}$  and R'.

However, this step must be repeated at each iteration This method sometimes gives spectacular result! However, often it gives irregular solutions.

### **Stopping Rules**

#### • $\chi^2$ variation

$$\chi_k^2 = ||y - R * \mu_k||^2 = \sum_{i}^{M \times N} \frac{(y_i - \sum_{ij} R_{ij} \mu_{j(k)})^2}{\sum_{ij} R_{ij} \mu_{j(k)}}$$
$$\frac{\chi_k^2 - \chi_{k-1}^2}{\chi_{k-1}^2} < 10^{-6}$$

If the regularization is good, one has  $\chi_k^2 \simeq \text{DoF}$ .

• Signal to noise ratio (usually measured in decibel)

$$SNR = 10 \log_{10} \left[ \frac{\sum_{i} (\mu_i - y_i)^2}{\sum_{i} (\mu_i - \mu_{true \ i})^2} \right]$$

where  $\mu_{true}$  is the true image. This quantity is used in the MC simulations during when the true image is known.

• convergence of the solution

$$\frac{||\mu_k - \mu_{k-1}||^2}{||\mu_{k-1}||^2} < 10^{-6}$$

Unfolding techniques

Statistica III